Rigorous free fermion entanglement renormalization from wavelets

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w/ Jutho Haegeman, Brian Swingle, Jordan Cotler, Glen Evenbly, and Volkher Scholz. See arXiv:1707.06243.

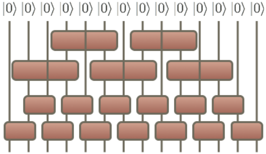
Tensor network states

Efficient variational classes for many-body quantum systems:

$$|\Psi\rangle = \sum_{i_1,\ldots,i_n} \Psi_{i_1,\ldots,i_n} |i_1,\ldots,i_n\rangle$$

e.g.

matrix product states



MERA

- can have interpretation as quantum circuit

Useful theoretical formalism:

- geometrize entanglement structure: *generalized area law*
- bulk-boundary dualities: *lift physics to the virtual level*
- quantum phases, topological order, RG, holography, . . . $\rightsquigarrow$ other talks

Tensor networks and quantum field theories

- tensor networks are discrete and finite representations
- quantum field theories are infinite and defined in the continuum

Two successful approaches:

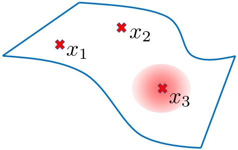
- *Lattice:* MPS, PEPS, MERA
- *Continuum:* cMPS, cMERA

How to measure goodness of approximation?
What does the tensor network describe?

Tensor networks for correlation functions

Given many-body system in state ρ and choice of operators $\{O_\alpha\}$, define correlation function:

$$C(\alpha_1, \cdots, \alpha_n) = \mathrm{tr}[\rho O_{\alpha_1} \cdots O_{\alpha_n}]$$



Goal: Design tensor network for correlation functions!

- unified perspective: system can be continuous, discreteness imposed by how we probe it
- tensor network for state sufficient — but not optimal
- in lattice models can recover state, but *only* for complete set of $\{O_\alpha\}$

Examples: Zaletel-Mong (MPS/q. Hall states), König-Scholz (MPS/CFTs), cf. quantum marginal problem

Our results

We construct tensor networks for free fermion systems:

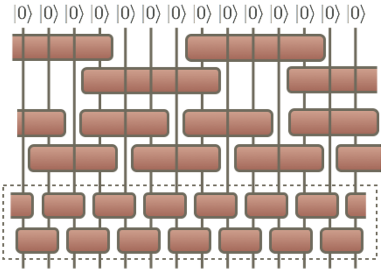
- 1D Dirac fermions on lattice & continuum
- Non-relativistic 2D fermions on lattice

Key features:

- Rigorous approximation of correlation functions
- Quantum circuits: MERA & branching MERA (Fermi surface)
- Explicit circuit construction, no variational optimization required

Continuum Dirac fermions $\rightsquigarrow$ upcoming paper w/ Scholz & Swingle

MERA: multi-scale entanglement renormalization ansatz (Vidal)



↓ quantum circuit that prepares state from $|0\rangle^{\otimes N}$

↑ entanglement renormalization, organize q. information by scale

- variational class for critical systems in 1D
- conjectured relationship to holography (Swingle)

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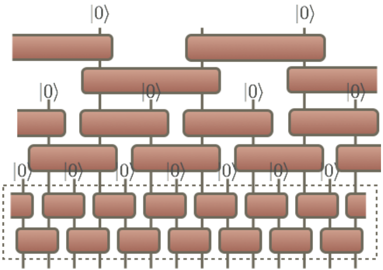

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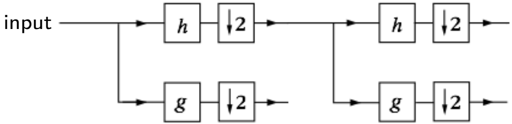
MERA and wavelets

Wavelet transforms organize classical information by scale:



- resolves discrete input signal in $\ell^2(\mathbb{Z})$ into different scales
- defined by low-pass ('scaling') filter h and high-pass ('wavelet') filter g

Key fact: Second quantizing 1D wavelet transform $\rightsquigarrow$ MERA circuit!

- in fact, obtain 'holographic' mapping (Qi)
- length of filter $\sim$ depth of layers (Evenbly-White)

Task: To produce free fermion ground state, design wavelet transform adapted to positive/negative energy modes.

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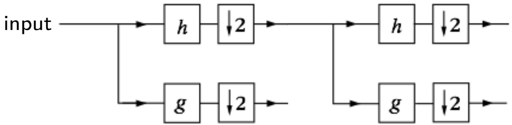

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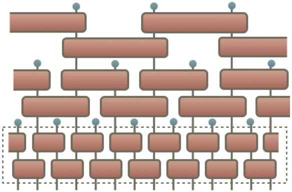

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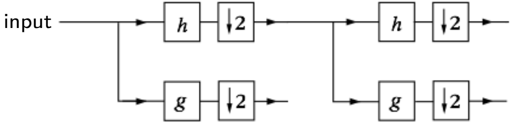

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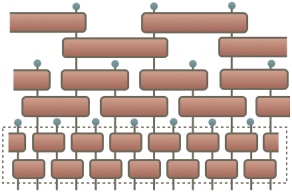

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1D Dirac fermions – Lattice model

Massless Dirac fermions on 1D lattice (Kogut-Susskind):

$$H_{1\mathrm{D}} = -\sum_n b^\dagger_{1,n} b_{2,n} - b^\dagger_{2,n} b_{1,n+1} + b^\dagger_{2,n} b_{1,n} - b^\dagger_{1,n+1} b_{2,n}$$

$$= \int_{-\pi}^{\pi} \frac{\mathrm{d}k}{2\pi} \begin{bmatrix} b_1(k) \\ b_2(k) \end{bmatrix}^\dagger \begin{bmatrix} 0 & \mathrm{e}^{-\mathrm{i}k} - 1 \\ \mathrm{e}^{\mathrm{i}k} - 1 & 0 \end{bmatrix} \begin{bmatrix} b_1(k) \\ b_2(k) \end{bmatrix}.$$

Diagonalize:

$$u(k) = \begin{bmatrix} 1 & 0 \\ 0 & -\mathrm{i}\,\mathrm{sign}(k)\mathrm{e}^{\mathrm{i}k/2} \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}, \quad u^\dagger h u = \begin{bmatrix} E_-(k) & 0 \\ 0 & E_+(k) \end{bmatrix}$$

- Fourier trafo highly *nonlocal*. But can choose *any* basis of Fermi sea!
- want pairs of modes related by $-\mathrm{i}\,\mathrm{sign}(k)\mathrm{e}^{\mathrm{i}k/2}$.

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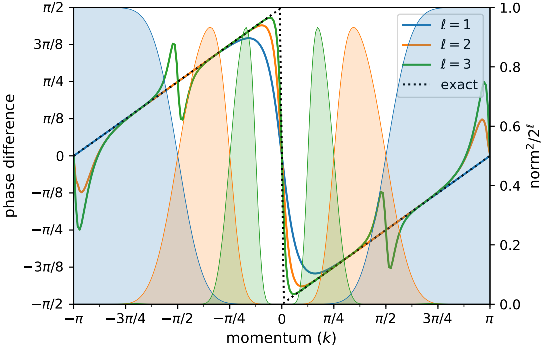
1D Dirac fermions – Wavelets

Equivalent: Pair of wavelet transforms such that high-pass filters are related by $-\mathrm{i}\,\mathrm{sign}(k)\mathrm{e}^{\mathrm{i}k/2}$.

- studied in signal processing, motivated by *translation-invariance*
- impossible with finite filters, but possible to arbitrary accuracy (Selesnick)



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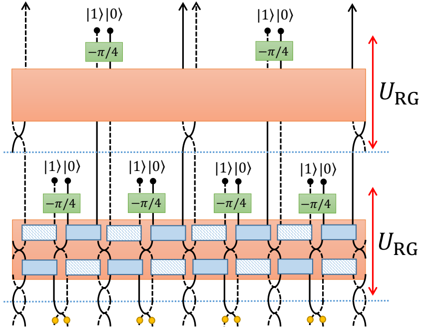


Parameters:

- $\mathcal{L}$ – number of layers
- ε – accuracy of phase relation of high-pass filters
- W – "size" of filters

Consider correlation function of N creation and annihilation operators

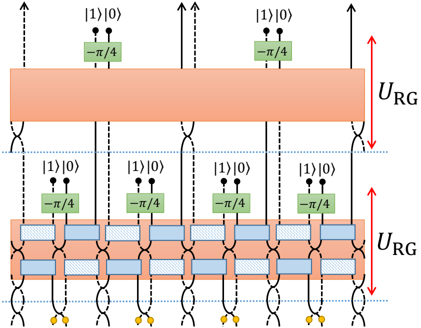
$$C(\{f_i\}) := \langle b^\dagger_{j_1}(f_1)\cdots b^\dagger_{j_N}(f_N)\, b_{j_{N+1}}(f_{N+1})\cdots b_{j_{2N}}(f_{2N})\rangle$$

supported on S lattice sites.

Theorem (simplified)

$$\left|C(\{f_i\})_{\text{exact}} - C(\{f_i\})_{\text{MERA}}\right| \lesssim \sqrt{SN}W\max\{2^{-\mathcal{L}/4}, \varepsilon\}$$

1D Dirac fermions – MERA



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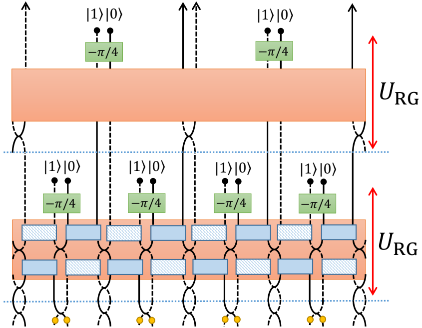

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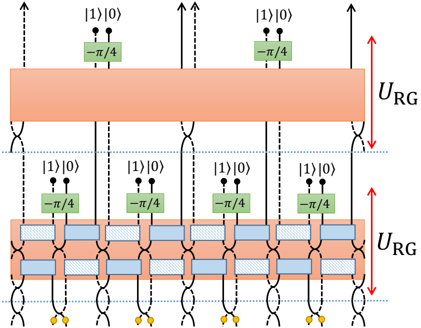

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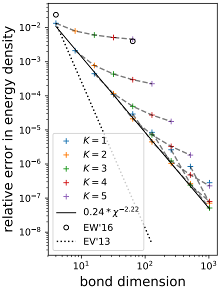
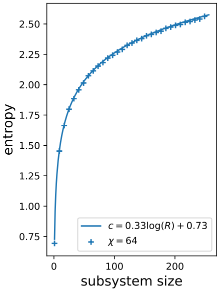
1D Dirac fermions – Numerics

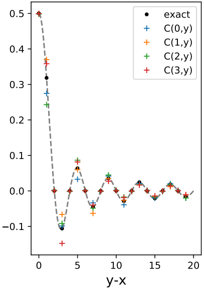
Energy error



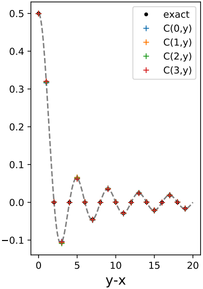
Entropy scaling



Green function $C(x, y) = \langle a_x^\dagger a_y \rangle$



bond dimension 2^2



bond dimension 2^6

1D Dirac fermions – Continuum

Massless Dirac fermions in (1+1)d:

$$\begin{bmatrix} \mathrm{i}(\partial_t + \partial_x) & 0 \\ 0 & \mathrm{i}(\partial_t - \partial_x) \end{bmatrix} \begin{bmatrix} \psi_+ \\ \psi_- \end{bmatrix} = 0$$

- need to produce modes $\psi_\pm(x)$ supported in $k < 0$ / $k > 0$

Natural construction: 'Continuum limit' of inverse wavelet transform!

- for *pair* of transforms as before: outputs $\psi_{1/2}$ related by $\mathrm{i}\,\mathrm{sign}(k)$
  $\rightsquigarrow \psi_\pm = \psi_1 \pm \mathrm{i}\psi_2$

Result: Rigorous quantum circuits for a quantum field theory!

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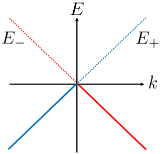

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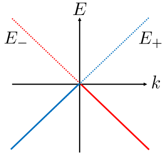

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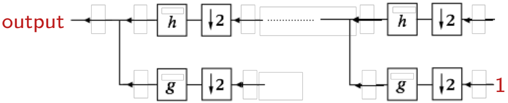

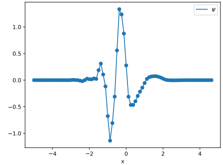

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Non-relativistic 2D fermions – Lattice model

$$H_{1D} \cong -\sum_n a_n^\dagger a_{n+1} + h.c.$$

Non-relativistic fermions hopping on 2D square lattice at half filling:

$$H_{2D} = -\sum_{m,n} a_{m,n}^\dagger a_{m+1,n} + a_{m,n}^\dagger a_{m,n+1} + h.c.$$

Fermi surface:

- Green function factorizes w.r.t. rotated axes
- violation of area law: $S(R) \sim R \log R$ (Wolf, Gioev-Klich, Swingle)

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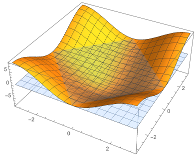

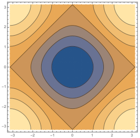

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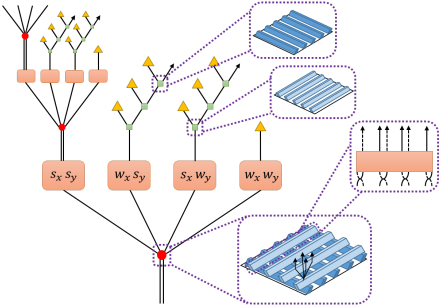
Non-relativistic 2D fermions – Branching MERA

Natural construction: Tensor product of wavelet transforms!

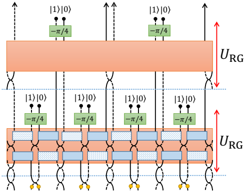
$$W\psi = \psi_s \oplus \psi_w \quad \rightsquigarrow \quad (W \otimes W)\psi = \psi_{ss} \oplus \psi_{ws} \oplus \psi_{sw} \oplus \psi_{ww}$$

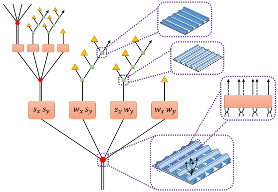
After second quantization, obtain variant of branching MERA (Evenbly-Vidal):



Similar approximation theorem holds.

Summary and outlook





Entanglement renormalization for free fermions:

- Rigorous approximation of correlation functions
- Explicit quantum circuits from wavelet transforms

Outlook:

- Massive theories, Dirac cones, beyond states at fixed times, ...
- Wess-Zumino-Witten CFTs (Scholz-Swingle-W.)
- Interacting theories? Starting point for variational optimization?

Thank you for your attention!

PhD & post-doc positions available @ University of Amsterdam / QuSoft