

Rigorous free fermion entanglement renormalization from wavelets

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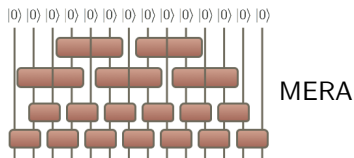
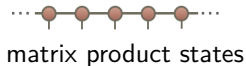
w/ Jutho Haegeman, Brian Swingle, Jordan Cotler, Glen Evenbly, and
Volkher Scholz. See [arXiv:1707.06243](https://arxiv.org/abs/1707.06243).

Tensor network states

Efficient variational classes for many-body quantum systems:

$$|\Psi\rangle = \sum_{i_1, \dots, i_n} \Psi_{i_1, \dots, i_n} |i_1, \dots, i_n\rangle$$

e.g.



- ▶ can have interpretation as **quantum circuit**

Useful theoretical formalism:

- ▶ geometrize **entanglement structure**: *generalized area law*
- ▶ bulk-boundary **dualities**: *lift physics to the virtual level*
- ▶ quantum phases, topological order, RG, holography, ... \rightsquigarrow other talks

Tensor networks and quantum field theories

- ▶ tensor networks are **discrete** and **finite** representations
- ▶ quantum field theories are **infinite** and defined in the **continuum**

Two successful approaches:

- ▶ *Lattice*: MPS, PEPS, MERA
- ▶ *Continuum*: cMPS, cMERA

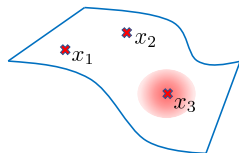
How to measure goodness of approximation?

What does the tensor network really capture?

Tensor networks for correlation functions

Given many-body system in state ρ and choice of operators $\{O_\alpha\}$, define **correlation function**:

$$C(\alpha_1, \dots, \alpha_n) = \text{tr}[\rho O_{\alpha_1} \cdots O_{\alpha_n}]$$



Goal: Design tensor network for correlation functions!

- ▶ unified perspective: system can be continuous, discreteness imposed by how we probe it
- ▶ tensor network for state sufficient — but not optimal
- ▶ in lattice models can recover state, but *only* for complete set of $\{O_\alpha\}$

Examples: Zaletel-Mong (MPS/q. Hall states), König-Scholz (MPS/CFTs),
cf. **quantum marginal problem**

Our results

We construct tensor networks for **free fermion systems**:

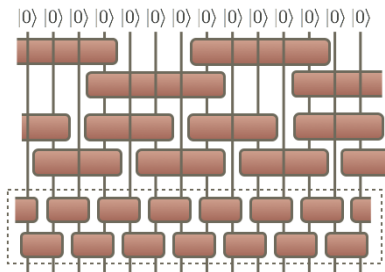
- ▶ 1D Dirac fermions on lattice & continuum
- ▶ non-relativistic 2D fermions on lattice (Fermi surface)

Key features:

- ▶ tensor networks that target **correlation functions**
- ▶ **rigorous** approximation guarantees
- ▶ **quantum circuits** that 'renormalize entanglement': (branching) MERA
- ▶ **explicit** circuit construction, no variational optimization required

Continuum results \rightsquigarrow upcoming paper w/ Scholz & Swingle

MERA: multi-scale entanglement renormalization ansatz (Vidal)



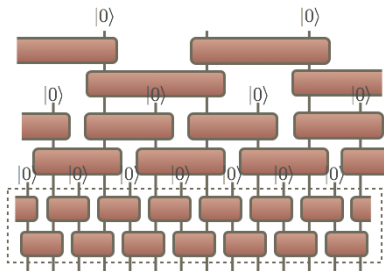
↓ local quantum circuit that prepares state from $|0\rangle^{\otimes N}$

↑ entanglement renormalization

↕ organize q. information by scale

- ▶ layers are short-depth quantum circuits (disentangle & coarse-grain)
- ▶ variational class for critical systems in 1D
- ▶ any MERA can be extended to a 'holographic' mapping
- ▶ reminiscent of holography (Swingle), starting point for tensor network models (HaPPY; Hayden-...-W.)

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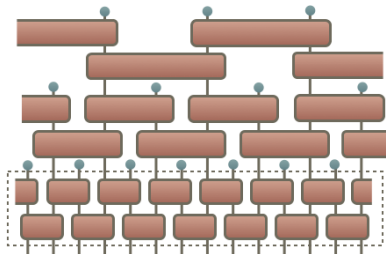
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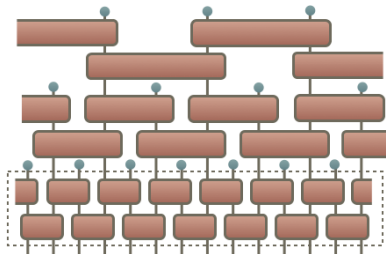
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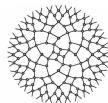


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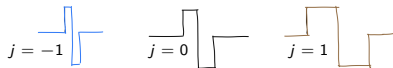


Wavelets

Wavelet transforms resolve **classical signal** into different scales, yet are local:

▶ multi-resolution analysis: $L^2(\mathbb{R}) = \bigoplus_j W_j$, spanned by $\psi(2^{-j}x - n)$

▶ ψ is called the *wavelet function*



▶ smooth & local functions need few W_j 's

Given signal at scale **up to j** , $V_j = \bigoplus_{j' \geq j} W_{j'}$, how to resolve it into scales?

- ▶ V_j spanned by $\phi(2^{-j}x - n)$,
- ▶ ϕ is known as *scaling function*

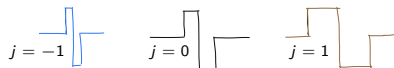
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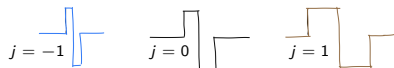
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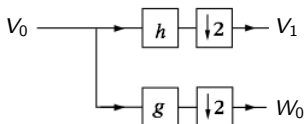
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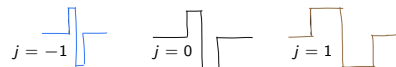


▶ defined by low-pass filter h and high-pass filter g

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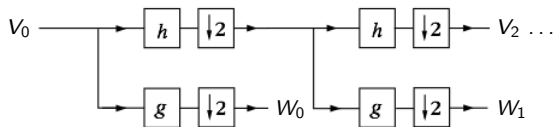
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$$V_0 = W_0 + V_1$$

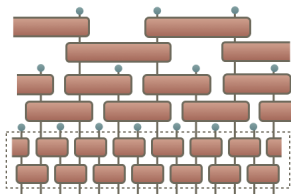
Discrete wavelet transform:



- ▶ defined by low-pass filter h and high-pass filter g
- ▶ locally resolves discrete input signal in $\ell^2(\mathbb{Z})$ into different scales

MERA and wavelets

Key fact: Second quantizing 1D wavelet transform \leadsto MERA circuit!

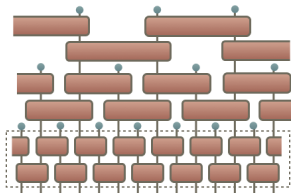


- ▶ in fact, obtain 'holographic' mapping (Q_i)
- ▶ length of classical filter \sim depth of quantum circuit (Evenbly-White)

Task: To produce free fermion ground state, design wavelet transform that targets positive/negative energy modes.

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1D Dirac fermions – Lattice model

Massless Dirac fermions on 1D lattice (Kogut-Susskind):

$$H_{1D} = - \sum_n b_{1,n}^\dagger b_{2,n} - b_{2,n}^\dagger b_{1,n+1} + b_{2,n}^\dagger b_{1,n} - b_{1,n+1}^\dagger b_{2,n}$$
$$= \int_{-\pi}^{\pi} \frac{dk}{2\pi} \begin{bmatrix} b_1(k) \\ b_2(k) \end{bmatrix}^\dagger \begin{bmatrix} 0 & e^{-ik} - 1 \\ e^{ik} - 1 & 0 \end{bmatrix} \begin{bmatrix} b_1(k) \\ b_2(k) \end{bmatrix}.$$

Diagonalize:

$$u(k) = \begin{bmatrix} 1 & 0 \\ 0 & -i \operatorname{sign}(k) e^{ik/2} \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}, \quad u^\dagger h u = \begin{bmatrix} E_-(k) & 0 \\ 0 & E_+(k) \end{bmatrix}$$

- ▶ Fourier trafo highly *nonlocal*. But can choose *any* basis of Fermi sea!
- ▶ want *pairs* of modes related by $-i \operatorname{sign}(k) e^{ik/2}$.

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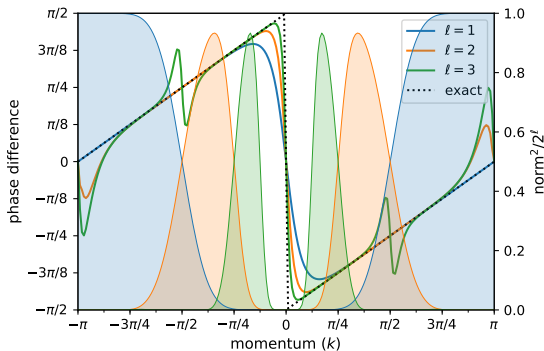
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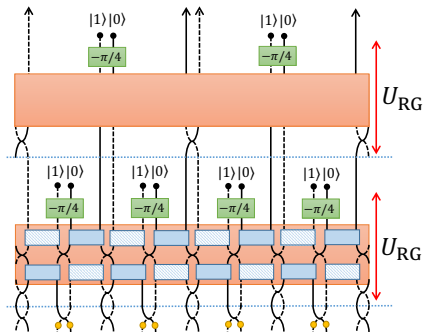
1D Dirac fermions – Wavelets

Task: Find pair of wavelet transforms such that **high-pass filters** are related by $-i \text{sign}(k)e^{ik/2}$.

- ▶ studied in signal processing, motivated by *translation-invariance* (!)
- ▶ impossible with finite filters, but possible to arbitrary accuracy (Selesnick)



1D Dirac fermions – MERA



Parameters:

- ▶ \mathcal{L} – number of layers
- ▶ ε – accuracy of phase relation of filters
- ▶ W – “size” of filters

Consider correlation function of N creation and annihilation operators

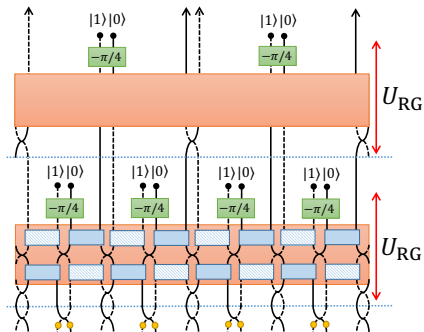
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supported on S lattice sites.

Theorem (simplified)

$$|C(\{f_i\})_{\text{exact}} - C(\{f_i\})_{\text{MERA}}| \lesssim \sqrt{SNW} \max\{2^{-\mathcal{L}/4}, \varepsilon\}$$

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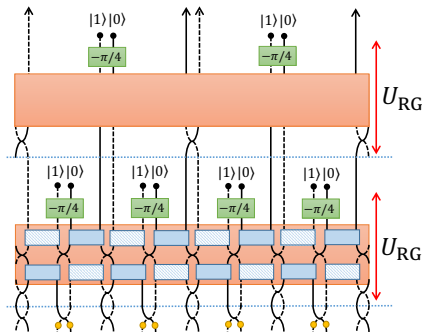
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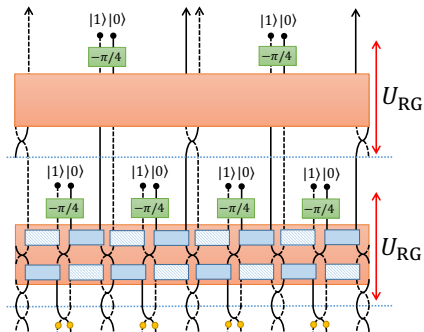
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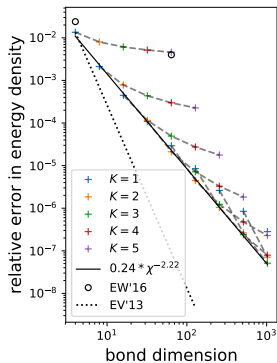
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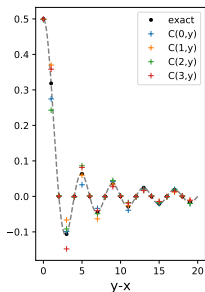
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1D Dirac fermions – Numerics

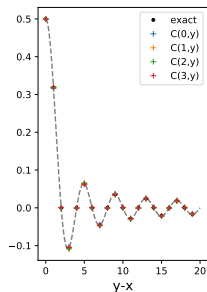
Energy error



Green function $C(x, y) = \langle a_x^\dagger a_y \rangle$



bond dimension 2^2



bond dimension 2^6

1D Dirac fermions – Continuum

Massless Dirac fermions in (1+1)d:

$$\begin{bmatrix} i(\partial_t + \partial_x) & 0 \\ 0 & i(\partial_t - \partial_x) \end{bmatrix} \begin{bmatrix} \psi_+ \\ \psi_- \end{bmatrix} = 0$$

- ▶ need to produce modes $\psi_{\pm}(x)$ supported in $k < 0$ / $k > 0$

Natural construction: ‘Continuum limit’ of inverse wavelet transform!

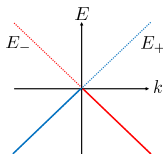
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Result: Rigorous quantum circuits for a quantum field theory!

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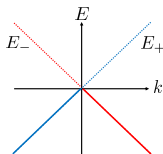
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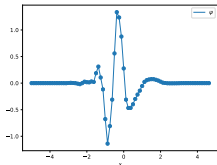
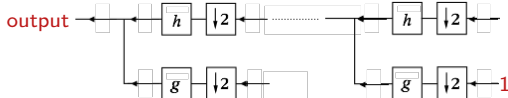
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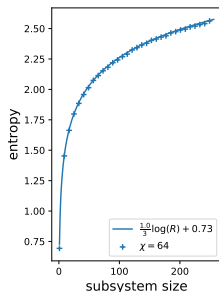


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1D Dirac fermions – Extracting conformal data

- ▶ **central charge:** $S(R) = \frac{c}{3} \log R + c'$
- ▶ **scaling dimensions:** find operators that coarse-grain to themselves (figures from Evenly-Vidal)



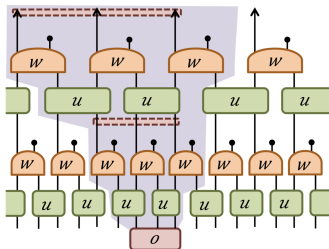
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Similarly: evaluate **quantum error correction** capabilities (cf. Kim-Kastoryano)

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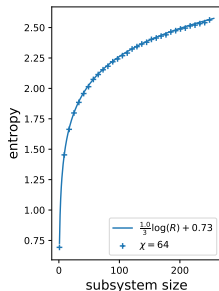
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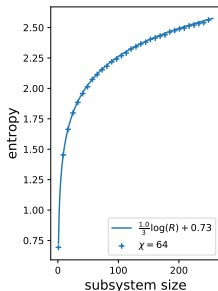
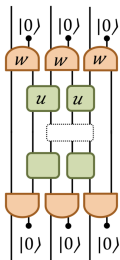
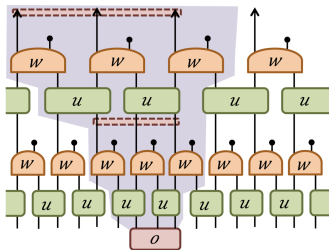
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$$H_{1D} \cong - \sum_n a_n^\dagger a_{n+1} + h.c.$$

Non-relativistic fermions hopping on 2D square lattice at half filling:

$$H_{2D} = - \sum_{m,n} a_{m,n}^\dagger a_{m+1,n} + a_{m,n}^\dagger a_{m,n+1} + h.c.$$

Fermi surface:

- ▶ violation of area law: $S(R) \sim R \log R$ (Wolf, Gioev-Klich, Swingle)
- ▶ Green function factorizes w.r.t. rotated axes

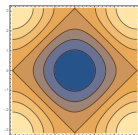
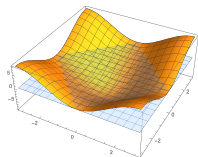
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Fermi surface:



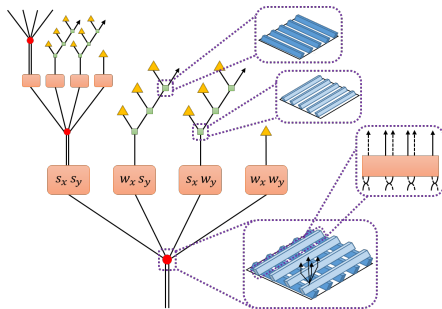
- ▶ violation of area law: $S(R) \sim R \log R$ (Wolf, Gioev-Klich, Swingle)
- ▶ Green function factorizes w.r.t. rotated axes

Non-relativistic 2D fermions – Branching MERA

Natural construction: **Tensor product** of wavelet transforms!

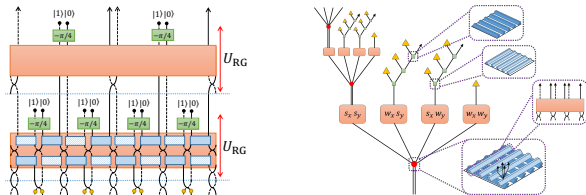
$$W\psi = \psi_s \oplus \psi_w \quad \rightsquigarrow \quad (W \otimes W)\psi = \psi_{ss} \oplus \psi_{ws} \oplus \psi_{sw} \oplus \psi_{ww}$$

After second quantization, obtain variant of **branching MERA** (Evenbly-Vidal):



Similar approximation theorem holds.

Summary and outlook



Entanglement renormalization for free fermions:

- ▶ Rigorous approximation of **correlation functions**
- ▶ Explicit quantum circuits from **wavelet** transforms

Outlook:

- ▶ Massive theories, Dirac cones, beyond states at fixed times, ...
- ▶ Wess-Zumino-Witten CFTs (w/ Scholz & Swingle); building block ...
- ▶ Interacting theories? Starting point for variational optimization?

Thank you for your attention!