

Entanglement renormalization and CFT: Quantum circuits for the Dirac field in 1+1 dimensions

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It from Qubit workshop, Kyoto, June 2019

with Freek Witteveen, Volkher Scholz, Brian Swingle ([1905.08821](#))

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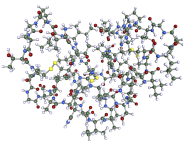
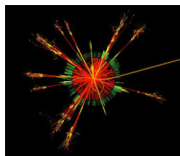


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Complexity of many-body quantum physics



Many-body states have **exponentially large** description:

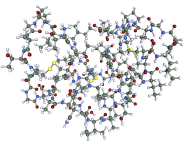
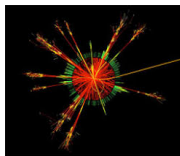
$$|\Psi\rangle = \sum_{i_1, \dots, i_n} \underbrace{\Psi_{i_1, \dots, i_n}}_{\text{exponentially large}} |i_1, \dots, i_n\rangle$$

In practice, entanglement **local** \leadsto compact description:

Start with local entangled pairs...

... and glue by applying **local transformations**:

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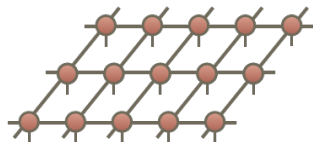


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Tensor networks in practice and theory

Tensor network: many-body state defined by contracting network of (local) tensors

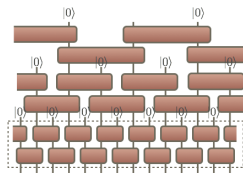


PEPS [Verstraete-Cirac]



matrix product state

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MERA [Vidal]

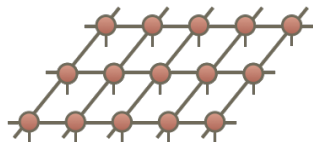
Numerical tool: **ansatz classes** for many-body states

- ▶ geometrize entanglement: *area (RT) laws*
- ▶ some are *quantum circuits*

Powerful **theoretical formalism** that provides 'dual' or 'holographic' descriptions of complex phenomena: topological order, ...

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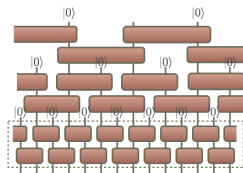


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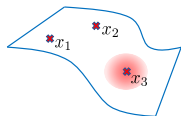
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Tensor networks and quantum field theory

Quantum field theories are defined in the **continuum**, while tensor networks are discrete and finitary. How to reconcile?

Two successful approaches:

- ▶ modify ansatz \leadsto continuum tensor networks
- ▶ relate discrete networks to **correlation functions** of continuum theory = unified perspective!



In either case. . .

What do tensor networks capture? Can we identify general construction principles? **Why do tensor networks work well?**

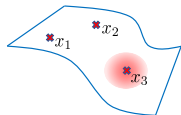
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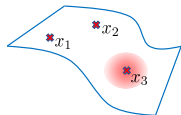
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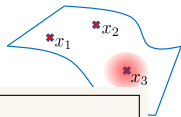
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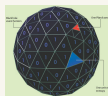
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More generally: **How to study QI in QFT?**

In ei ☺ notions like *subsystem*, *entropy*, *approximation*, etc. subtle

cf. pl ☺ exciting recent progress. large body of work in mathematical physics.



c-theorem from subadditivity, Bekenstein bound via relative entropy, renormalization vs QEC, approximate QEC, ...

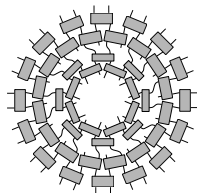
Our contributions

Result

We construct tensor networks for the **Dirac CFT** in 1+1 dimensions.

Key features:

- ▶ **explicit** construction – no variational optimization
- ▶ **rigorous** approximation of correlation functions
- ▶ quantum circuits that **renormalize entanglement**



We achieve this using tools from signal processing: multiresolution analysis and discrete/continuum duality from **wavelet theory**.

Also obtain sub-circuits for Majorana and Ising CFT. In prior work, we constructed (branching) MERA for critical free-fermion lattice models.

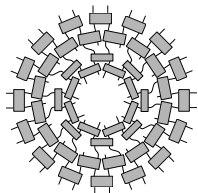
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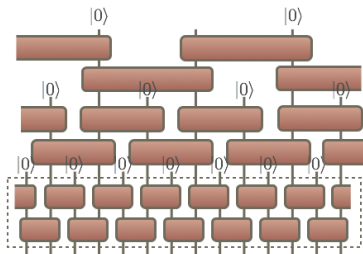


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Tensor network ansatz for **critical systems**: [Vidal]

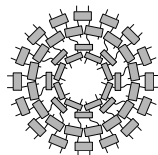


↓ local **quantum circuit** that prepares state from $|0\rangle^{\otimes N}$

↑ entanglement renormalization

↕ organize q. information by scale

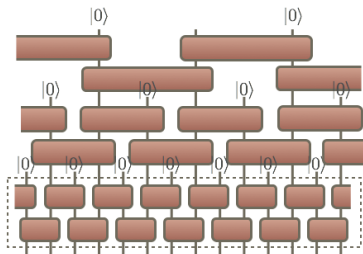
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- ▶ noise-resilient on quantum computer [Kim et al]
- ▶ reminiscent of **holography** [Swingle], starting point for **tensor network models** [Qi, HaPPY, Hayden...-W]



Important to understand design principles! Can we bridge numerics and toy models?

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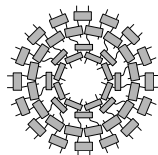


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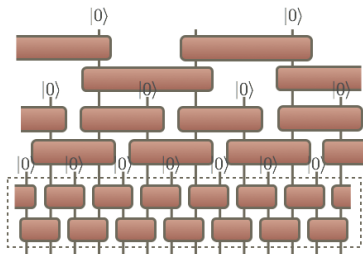
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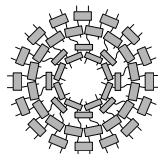


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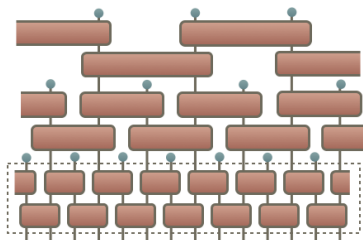
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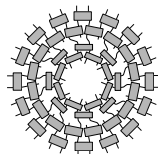


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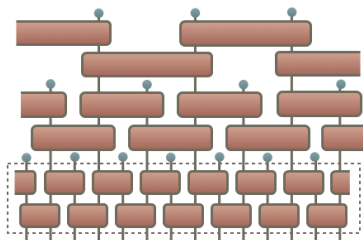
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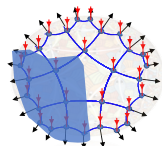


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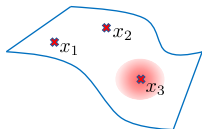
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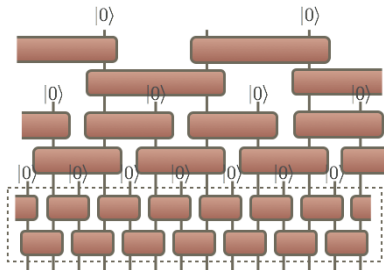
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Result: Entanglement renormalization for Dirac CFT

Massless Dirac fermion in 1+1d: $i\gamma^\mu \partial_\mu \psi = 0$



We construct MERAs that target vacuum **correlation functions** $C = \langle O_1 \cdots O_n \rangle$ of smeared observables.



Result (simplified)

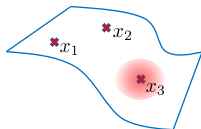
$$C_{\text{exact}} \approx C_{\text{MERA}}$$

Goodness depends on smearing, #layers, quality parameter. Comes with 'dictionary' for mapping observables. Symmetries approximately inherited!

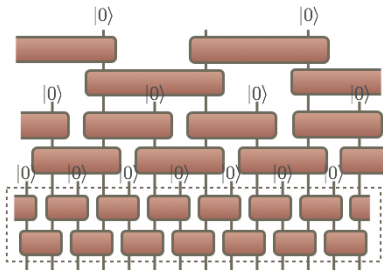
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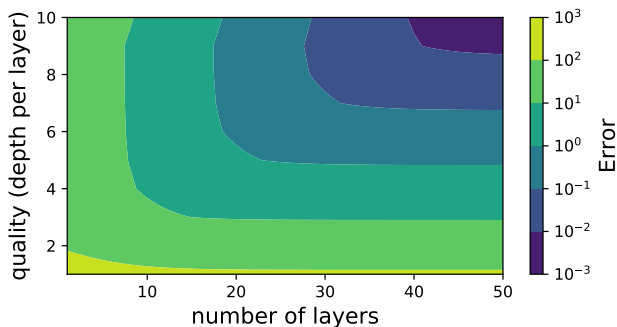
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Two-point functions: A-priori error

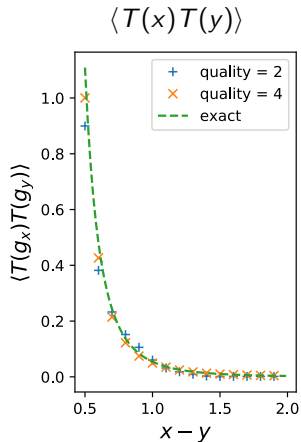
For different values of quality parameter and number of layers:



Similarly for higher-point functions.

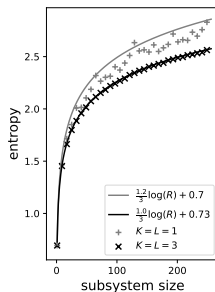
Two-point functions: Numerics

For different values of quality parameter and *large* number of layers:



Verifying conformal data

- ▶ **central charge:** $S(R) = \frac{c}{3} \log R + c'$
- ▶ usual procedure: identify fields by searching for operators that coarse-grain to themselves



→ diagonalize 'scaling superoperator' [Evenbly-Vidal]

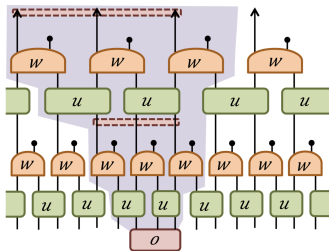
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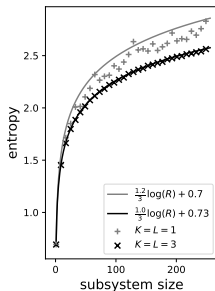
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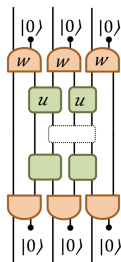
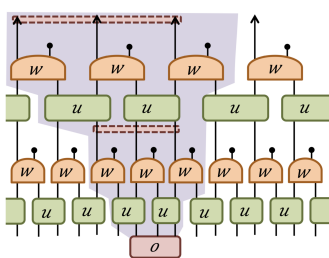
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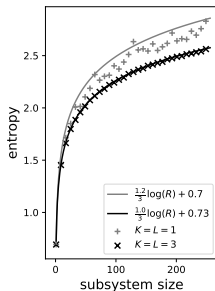
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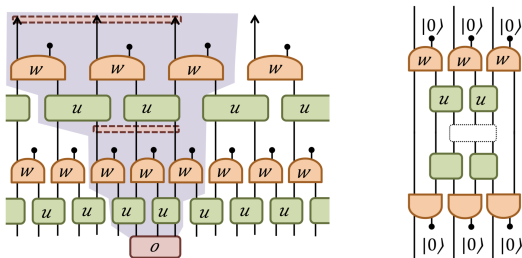
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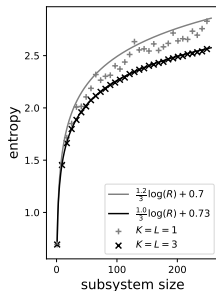
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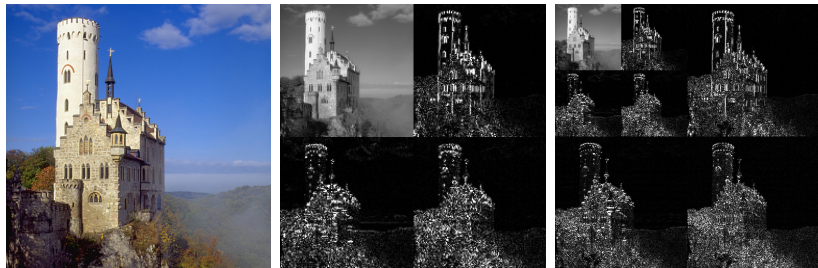
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Key technique: Entanglement renormalization using **wavelet theory**.

Tool from signal processing to resolve signal by **scale**:



Mathematically, basis transform built from scalings and translates of single localized 'wavelet'.

- ▶ second quantization: quantum circuit!

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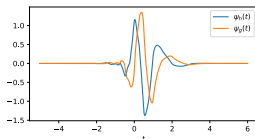
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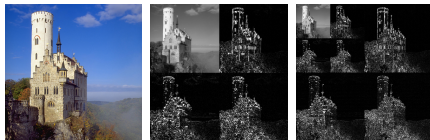


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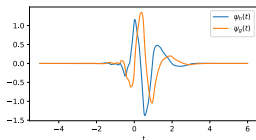
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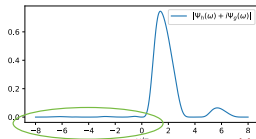


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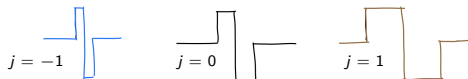


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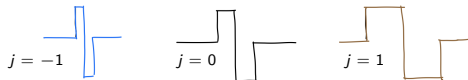
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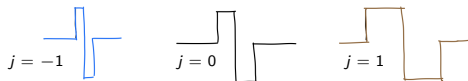
The equation shows a green wavelet packet on the left, followed by an equals sign, then a pink wavelet packet scaled by $\frac{1}{2}$, a plus sign, and another pink wavelet packet scaled by $\frac{1}{2}$. The pink packets are wider than the green packet.

Transform is implemented by circuit on *single-particle* Hilbert space:

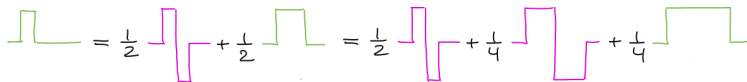
- ▶ Discrete circuit resolves continuous signal by scale!
- ▶ Second quantization yields Gaussian MERA layer. [Evenbly-White]

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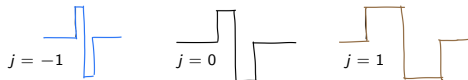


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In more detail: Wavelets and MERA

Wavelet bases are built from scalings & translates of single wave packet:



We can recursively resolve signal into different scales (multiresolution analysis):

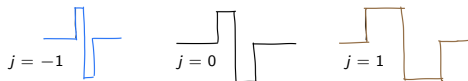
The equation shows a green wavelet packet on the left, followed by an equals sign, then a pink wavelet packet scaled by $\frac{1}{2}$, a plus sign, and another pink wavelet packet scaled by $\frac{1}{2}$. The pink packets are wider than the green packet.

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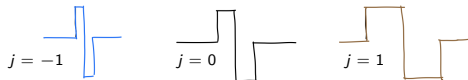
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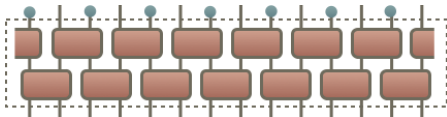


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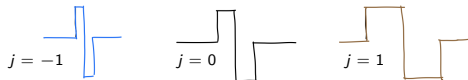


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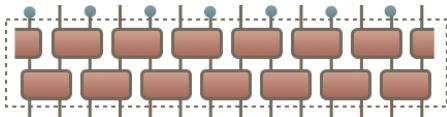
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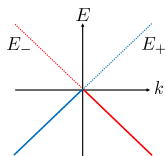
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Wavelets for the Dirac fermion

Massless Dirac equation in 1+1d:

$$\begin{bmatrix} \psi_1(x, t) \\ \psi_2(x, t) \end{bmatrix} = \begin{bmatrix} \chi_+(x - t) + \chi_-(x + t) \\ i\chi_+(x - t) - i\chi_-(x + t) \end{bmatrix}$$



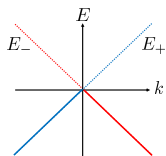
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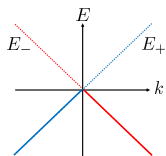
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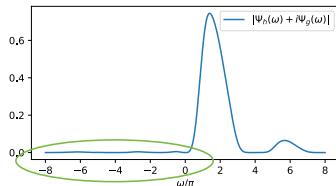
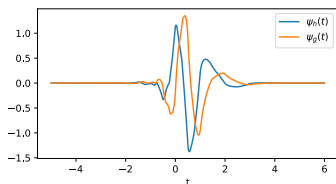
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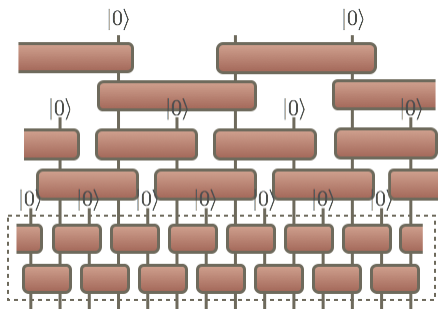
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Technical Result



Parameters:

- ▶ \mathcal{L} – number of layers
- ▶ ϵ – accuracy of wavelet pair
- ▶ Γ – support and smoothness of smearing functions

Consider **correlation function** with smeared fields & normal-ordered bilinears:

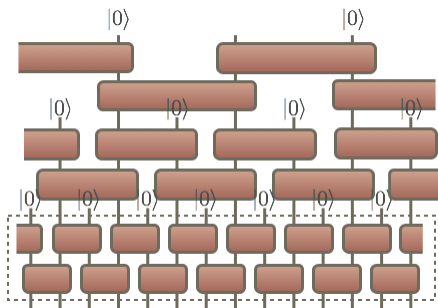
$$C := \langle \Psi^\dagger(f_1) \cdots \Psi(f_{2N}) O_1 \cdots O_M \rangle$$

Theorem (simplified)

$$|C_{\text{exact}} - C_{\text{MERA}}| \leq \Gamma \max\{2^{-\mathcal{L}/3}, \epsilon \log \frac{1}{\epsilon}\}$$

We provide dictionary for C_{MERA} (discretize smearing functions in scaling basis etc).

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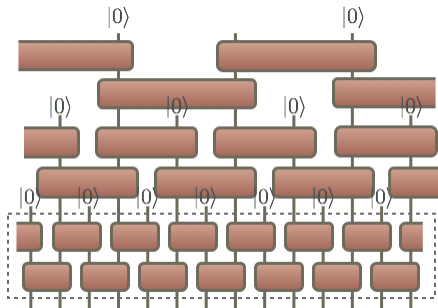
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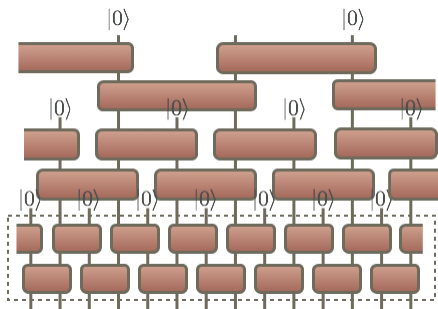
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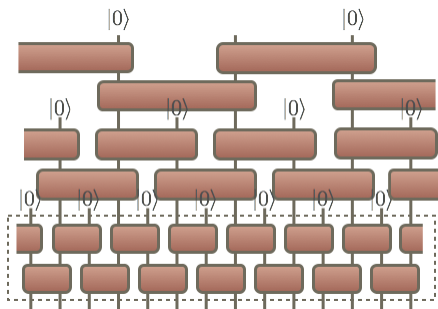
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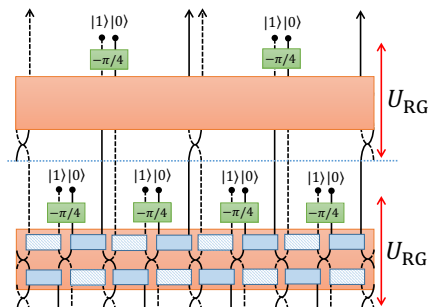
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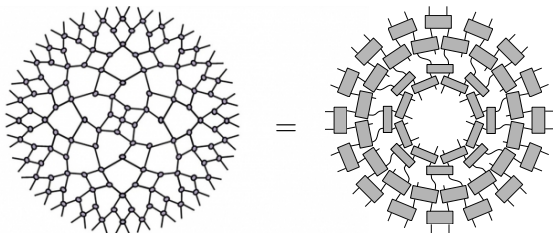
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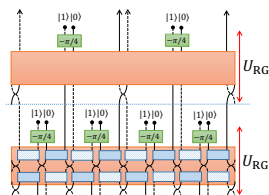
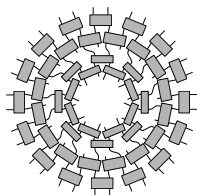
Dirac fermion on circle

Construction can be easily adapted to Dirac fermion on circle:



- ▶ finite number of layers once UV cut-off fixed
- ▶ systematic construction by (anti)periodizing wavelets

Summary and outlook



- ▶ entanglement renormalization **quantum circuits** for 1+1d Dirac CFT
- ▶ systematic construction with **rigorous guarantees**

Outlook:

- ▶ thermofield double, Dirac cones, ...
- ▶ **building block** for more interesting CFTs? **starting point** for perturbation theory or variational optimization?
- ▶ **lift** wavelet theory to quantum circuits! tensor network bootstrap?

Thank you for your attention!

How to build an approximate Hilbert pair

[Selesnick]

Wavelets are built from **filters** $g[n]$ that relate functions at different scale:

$$\phi_{j-1}(x) = \sum_{n \in \mathbb{Z}} g[n] \phi_j(x - 2^{-j} n)$$

Necessary and sufficient to obtain orthonormal basis (roughly speaking):

$$|G(\theta)|^2 + |G(\theta + \pi)|^2 = 2, \quad G(0) = \sqrt{2}$$

Wavelets are related by Hilbert transform iff filters related by **half-shift**:

$$G(\theta) = H(\theta) e^{-i\theta/2}$$

To achieve this, find explicit approximation

$$e^{-i\theta/2} \approx e^{-iL\theta} \frac{D(-\theta)}{D(\theta)}.$$

Then, $H(\theta) = Q(\theta)D(\theta)$ and $G(\theta) = Q(\theta)e^{-iL\theta}D(-\theta)$ are approximately related by half-shift for any choice of $Q(\theta)$.

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When put on lattice, massless Dirac fermion becomes: (Kogut-Susskind)

$$H_{1D} \cong - \sum_n a_n^\dagger a_{n+1} + h.c.$$

Non-relativistic fermions hopping on 2D square lattice at half filling:

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Fermi surface:

- ▶ violation of area law: $S(R) \sim R \log R$ (Wolf, Gioev-Klich, Swingle)
- ▶ Green's function factorizes w.r.t. rotated axes

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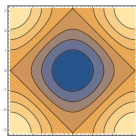
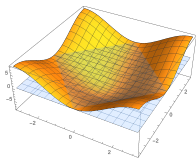
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Natural construction – perform wavelet transforms in both directions:

$$W\psi = \psi_{\text{low}} \oplus \psi_{\text{high}} \quad \rightsquigarrow \quad (W \otimes W)\psi = \psi_{ll} \oplus \psi_{lh} \oplus \psi_{hl} \oplus \psi_{hh}$$

After second quantization, obtain variant of **branching MERA** (Evenbly-Vidal):

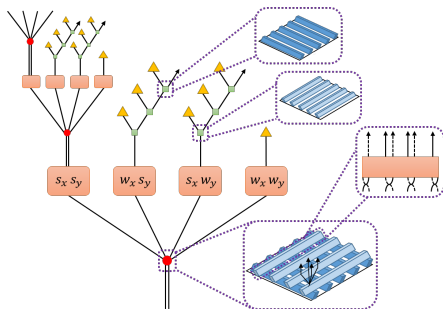
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