

Quantum circuits for the Dirac field in 1+1 dimensions

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Tensor Networks Workshop, AEI, March 2019

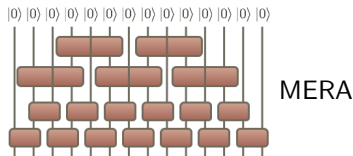
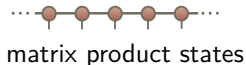
based on [arXiv:1905.08821](#) (with Witteveen, Scholz, Swingle) and [arXiv:1707.06243](#) (with Haegeman, Swingle, Cotler, Evenbly, Scholz)



Tensor networks

$$|\psi\rangle = \sum_{i_1, \dots, i_n} \boxed{\psi_{i_1, \dots, i_n}} |i_1, \dots, i_n\rangle$$

Efficient variational classes for many-body quantum states:



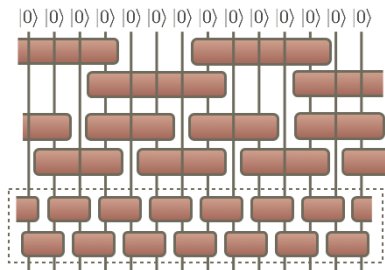
- ▶ can have interpretation as **quantum circuit**

Useful theoretical formalism:

- ▶ geometrize **entanglement structure**: *generalized area law*
- ▶ bulk-boundary **dualities**: *lift physics to the virtual level*
- ▶ quantum phases, topological order, RG, holography, ...

MERA

multi-scale entanglement renormalization ansatz (Vidal)



↓ local **quantum circuit** that prepares state from $|0\rangle^{\otimes N}$

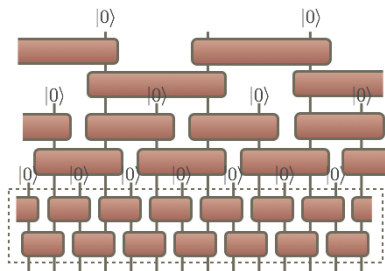
↑ entanglement renormalization

↕ organize q. information by scale

- ▶ self-similar layers that are short-depth quantum circuits
- ▶ variational class for **critical systems** in 1D
- ▶ interpretation: disentangle & coarse-grain
- ▶ network arises from **tensor network renormalization**:

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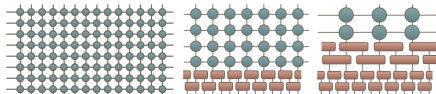


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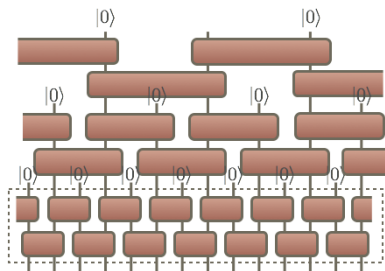
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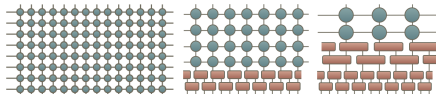


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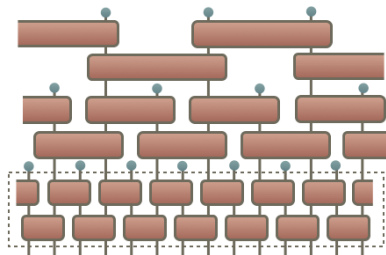
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MERA and holography



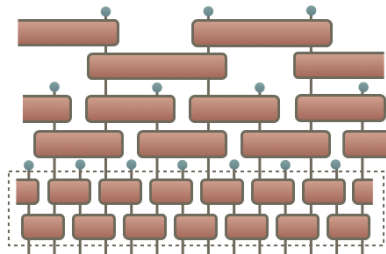
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- ▶ can always extend to ‘holographic’ mapping
- ▶ hyperbolic geometry (Swingle)
- ▶ starting point for **tensor network models** of holography (HaPPY; Hayden-...-W.)
- ▶ **quantum error correction** property = **noise-resilience** on QC (Kim et al)
 \leadsto important to understand design principles

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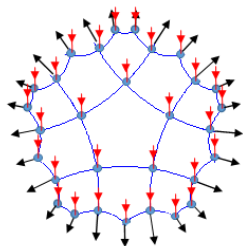


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Tensor networks and quantum field theories

Tensor networks are **discrete** and **finitary** representations, while QFTs are **infinite** and defined in the **continuum**.

Two successful approaches:

- ▶ *continuum* (cMPS, cMERA, ...)
- ▶ *lattice* (MPS, PEPS, MERA, ...)

Questions:

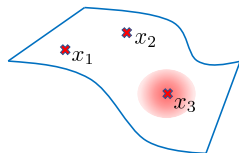
- ▶ **what do tensor networks capture?**
- ▶ how to measure goodness of approximation?
- ▶ can we give rigorous construction principles?
- ▶ **why do tensor networks work well?**

cf. plethora of results on gapped 1D lattice systems in QIT/cond-mat

Tensor networks for correlation functions

Given many-body system in state ρ and choice of operators $\{O_\alpha\}$, define **correlation function**:

$$C(\alpha_1, \dots, \alpha_n) = \text{tr}[\rho O_{\alpha_1} \cdots O_{\alpha_n}]$$



Goal: Design tensor network for correlation functions!

- ▶ unified perspective: system can be continuous – discreteness imposed by how we probe it
- ▶ tensor network for ρ sufficient (if possible), but likely suboptimal

Examples: Zaletel-Mong (MPS/q. Hall states), **König-Scholz** (MPS/CFTs),
cf. **quantum marginal problem**

Our results

We construct tensor networks for free fermion theories:

- ▶ 1D Dirac fermion in continuum & lattice
- ▶ non-relativistic 2D fermions on lattice (Fermi surface)

Key features:

- ▶ tensor networks that target correlation functions
- ▶ rigorous approximation guarantees
- ▶ entanglement renormalization quantum circuits: (branching) MERA
- ▶ explicit construction, no variational optimization

We achieve this using tools from signal processing: wavelet theory.

1D Dirac fermion – Lattice result

Fermions hopping on infinite 1D lattice at half filling:

$$H_{1D} = - \sum_n a_n^\dagger a_{n+1} + h.c.$$

- ▶ equivalent to ‘staggered’ massless Dirac fermions (Kogut-Susskind)
- ▶ easily solved using Fourier transform – but *not* using local q. circuit!

We construct MERA networks
that target **correlation functions**:

$$C(\{f_i\}) := \langle a^\dagger(f_1) \cdots a(f_{2N}) \rangle$$

Result (simplified)

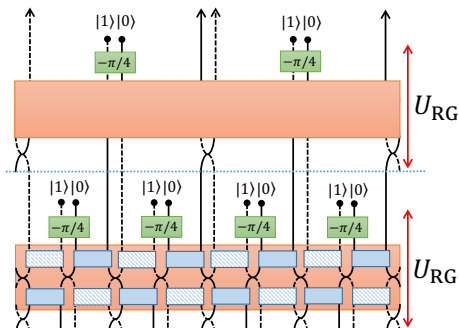
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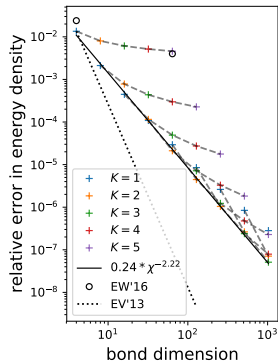
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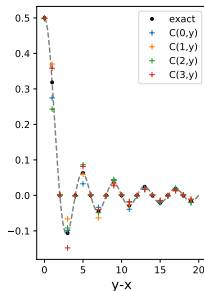
1D Dirac fermion – Numerics

Energy error

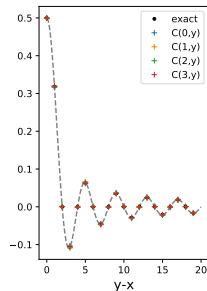


bond dimension = $2^{\text{circuit depth}}$

Green function $C(x, y) = \langle a_x^\dagger a_y \rangle$



bond dimension 2^2

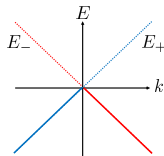


bond dimension 2^6

1D Dirac fermion – Continuum result

Massless Dirac fermion in 1+1d:

$$i\gamma^\mu \partial_\mu \psi = 0$$



We construct circuits that target **vacuum correlation functions** in Dirac CFT:

$$C(\{f_i, A_j\}) := \langle \Psi^\dagger(f_1) \cdots \Psi(f_{2N}) A_1 \cdots A_M \rangle$$

$\Psi(f_i)$ smeared fields, A_j normal-ordered bilinears (e.g. smeared T , L_n)

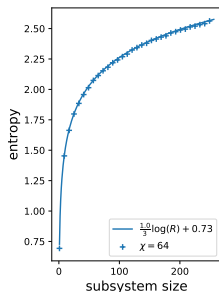
Result (simplified)

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*Rigorous quantum circuit
approximation for a QFT!*

1D Dirac fermion – Verifying conformal data

- ▶ **central charge:** $S(R) = \frac{c}{3} \log R + c'$
- ▶ usual procedure: identify fields by searching for operators that coarse-grain to themselves

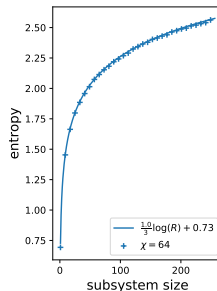
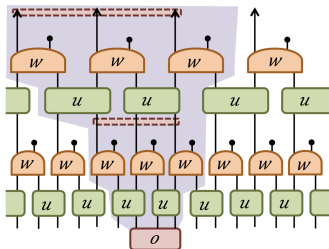


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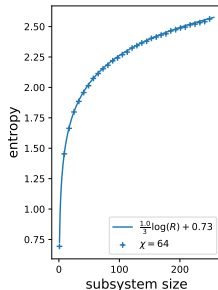
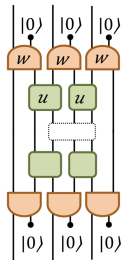
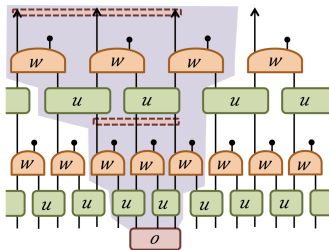


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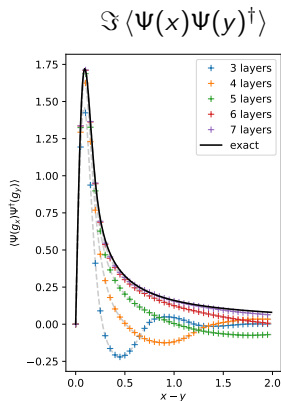
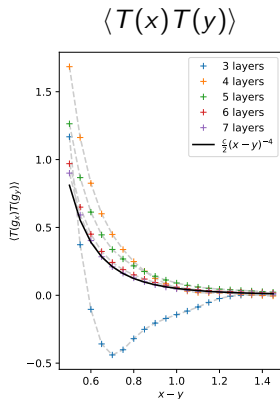


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1D Dirac fermion – Numerics

Two-point functions:

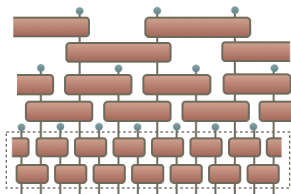


Similarly: OPE coefficients.

How to construct a free-fermion (= Gaussian) MERA?

Free-fermion ground states are Fermi seas filled with negative energy modes of single-particle Hamiltonian. This begs the question:

*How to perform entanglement renormalization on the single-particle level?
Is there a single-particle variant of MERA?*

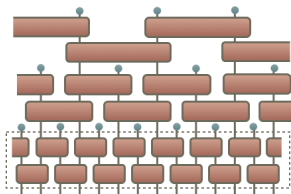


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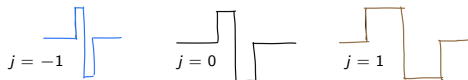
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Wavelets and renormalization

Fourier basis resolves signal into scales, but is completely nonlocal.
In contrast, can also generate basis by scalings and translates of single localized wave packet – a **wavelet**:



Then we can recursively resolve signal into different scales:

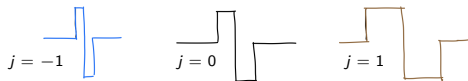
where

W_j = span of wavelets at scale j

V_j = signals at scale up to $j = W_j \oplus W_{j+1} \oplus \dots$

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$$\begin{aligned} V_0 &= \frac{1}{2} W_0 + \frac{1}{2} V_1 \\ V_0 &= \frac{1}{2} W_0 + \frac{1}{4} W_1 + \frac{1}{4} V_2 \end{aligned}$$

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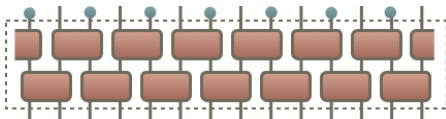
Wavelets and MERA

The basis transformation

$$V_j \rightarrow W_j \oplus W_{j+1}$$

$$\begin{array}{c} \text{[Step Function]} \\ V_0 \end{array} = \frac{1}{2} \begin{array}{c} \text{[Step Function]} \\ W_0 \end{array} + \frac{1}{2} \begin{array}{c} \text{[Step Function]} \\ V_1 \end{array}$$

is implemented by a *classical circuit* acting on *single-particle* Hilbert space:



Second quantization yields *layer of a Gaussian MERA!*

- ▶ in fact, obtain 'holographic' mapping (Qi)
- ▶ depth of classical circuit = depth of quantum circuit (Evenbly-White)

Upshot: To construct free-fermion ground state, *design wavelet transform* that targets positive/negative energy modes.

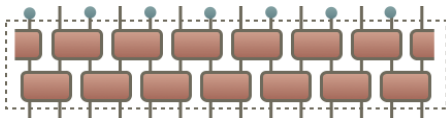
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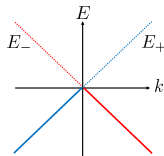
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1D Dirac fermion – Vacuum state

Massless Dirac equation in 1+1d:

$$i\gamma^\mu \partial_\mu \psi = 0$$



Negative energy modes:

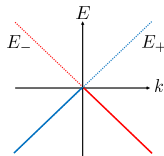
- ▶ χ_\pm supported on $k < 0$ / $k > 0$
- ▶ $\psi_{1,2}$ related by $-i \text{sign}(k)$ at $t = 0$ (*Hilbert transform*)
- ▶ can choose *any* basis of Fermi sea...

Goal: Design pair of wavelets related by Hilbert transform!

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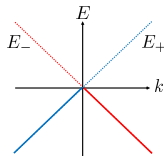
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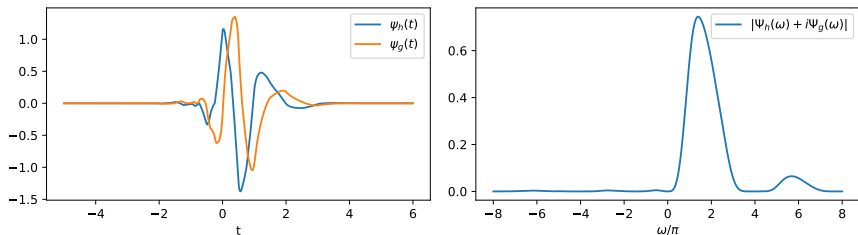
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1D Dirac fermion – Hilbert wavelet pairs

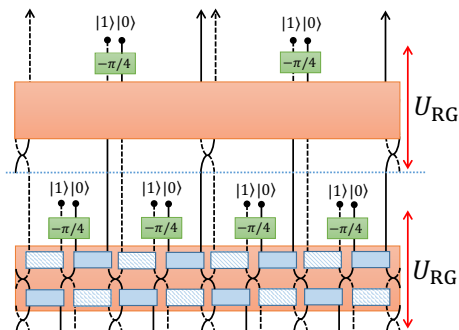
Such wavelet pairs have been studied in the signal processing community:

- ▶ motivated by *directionality* and *shift-invariance* (!)
- ▶ impossible exactly with local circuit, but possible to arbitrary accuracy (Selesnick)



After second quantizing and careful analysis, obtain tensor network with rigorous approximation guarantee. . .

1D Dirac fermion – Result



Parameters:

- ▶ \mathcal{L} – number of layers
- ▶ ε – accuracy of Hilbert pair
- ▶ Γ – support and smoothness of smearing functions

Consider **correlation function** with smeared fields & normal-ordered bilinears:

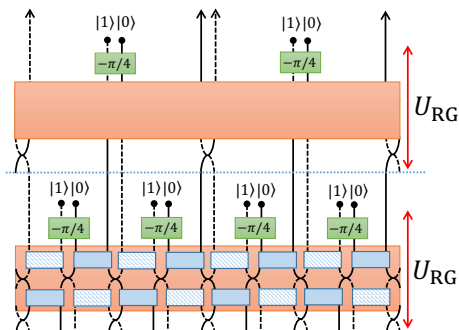
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Result (simplified)

$$|C_{\text{exact}} - C_{\text{MERA}}| \leq \Gamma \max\{2^{-\mathcal{L}/4}, \varepsilon \log \frac{1}{\varepsilon}\}$$

In particular, all conformal symmetries approximately inherited.

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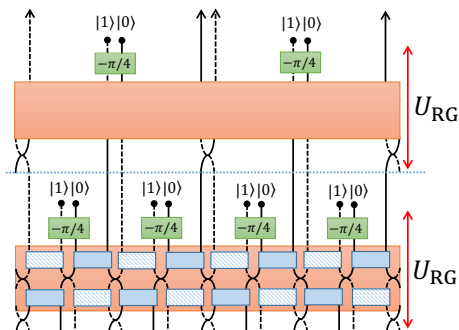
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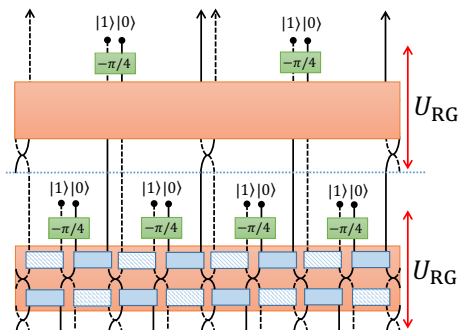
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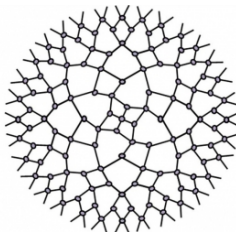
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In particular, all conformal symmetries approximately inherited.

1D Dirac fermion – Circle

Construction also works for Dirac fermion on circle:



- ▶ finite number of layers once UV cut-off fixed
- ▶ systematic construction by (anti)periodizing wavelets
- ▶ only top layers change – wavelet modes start ‘wrapping around’

Non-relativistic 2D fermions – Lattice model

$$H_{1D} \cong - \sum_n a_n^\dagger a_{n+1} + h.c.$$

Non-relativistic fermions hopping on 2D square lattice at half filling:

$$H_{2D} = - \sum_{m,n} a_{m,n}^\dagger a_{m+1,n} + a_{m,n}^\dagger a_{m,n+1} + h.c.$$

Fermi surface:

- ▶ violation of area law: $S(R) \sim R \log R$ (Wolf, Gioev-Klich, Swingle)
- ▶ Green's function factorizes w.r.t. rotated axes

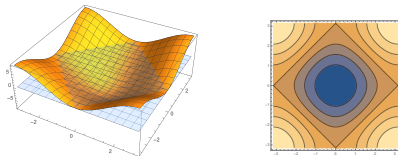
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Non-relativistic 2D fermions – Branching MERA

Natural construction – perform wavelet transforms in both directions:

$$W\psi = \psi_{\text{low}} \oplus \psi_{\text{high}} \quad \leadsto \quad (W \otimes W)\psi = \psi_{ll} \oplus \psi_{lh} \oplus \psi_{hl} \oplus \psi_{hh}$$

After second quantization, obtain variant of **branching MERA** (Evenbly-Vidal):

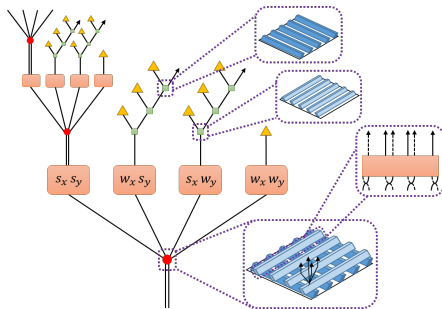
Similar approximation theorem holds.

Non-relativistic 2D fermions – Branching MERA

Natural construction – perform wavelet transforms in both directions:

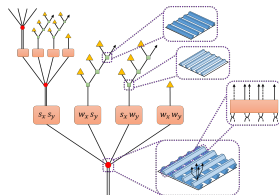
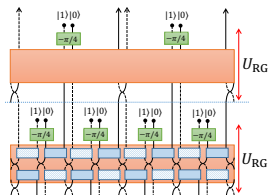
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After second quantization, obtain variant of **branching MERA** (Evenbly-Vidal):



Similar approximation theorem holds.

Summary and outlook



- ▶ entanglement renormalization **quantum circuits** for free fermions
- ▶ explicit construction with **rigorous guarantees** (lattice + continuum)

Outlook:

- ▶ thermofield double, massive theories, Dirac cones, ...
- ▶ **building block** for more interesting CFTs? **starting point** for perturbation theory or variational optimization?

Thank you for your attention!