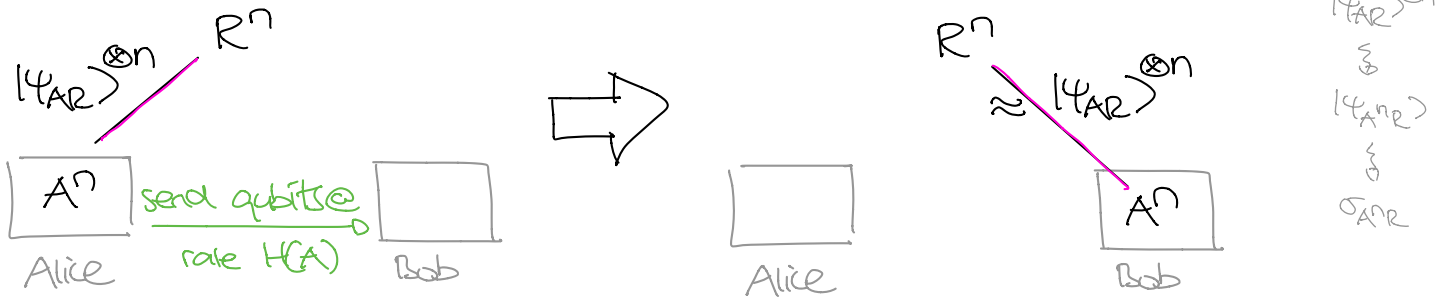


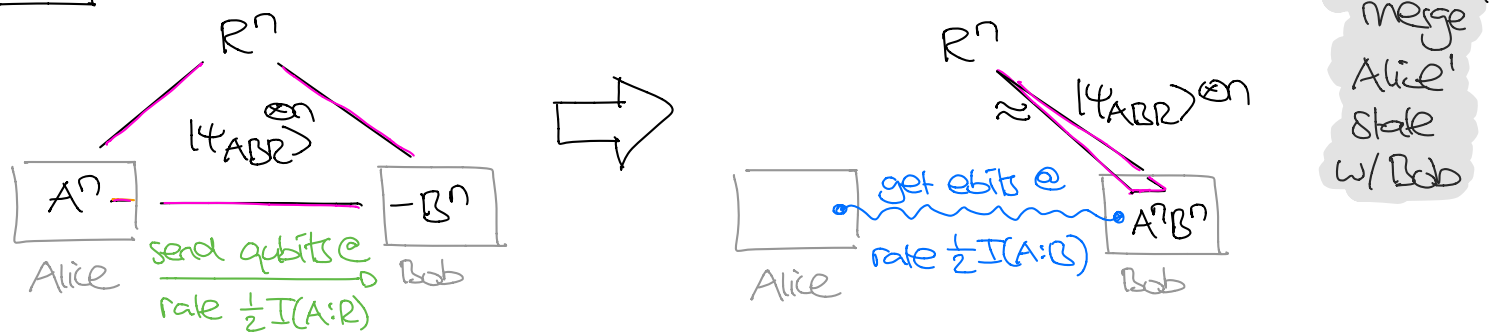
Quantum state merging

Recall Schumacher Compression: Given $|\psi_{AR}\rangle$, want



What if Bob already has part of state? Rate $H(A)$ is still possible! But...

GOAL:



(Other variants: state splitting (Alice keeps part of state) or combination)

* No B (or ψ_{AR} pure): Schumacher compression at rate $\frac{1}{2} I(A:R) = H(A)$ ✓
 in general: $\frac{1}{2} I(A:R) \leq H(A)$ and may even get entanglement...

* No R (or ψ_{AB} pure): Entanglement distillation at rate $\frac{1}{2} I(A:B) = H(B)$ ✓
 w/o communication

in general? given $|\psi_{AB}\rangle^{\otimes n}$, can distill ebits at rate $\frac{1}{2} I(A:B) - \frac{1}{2} I(A:R)$
 using LOCC
 $= H(B) - H(AB)$ if ≥ 0

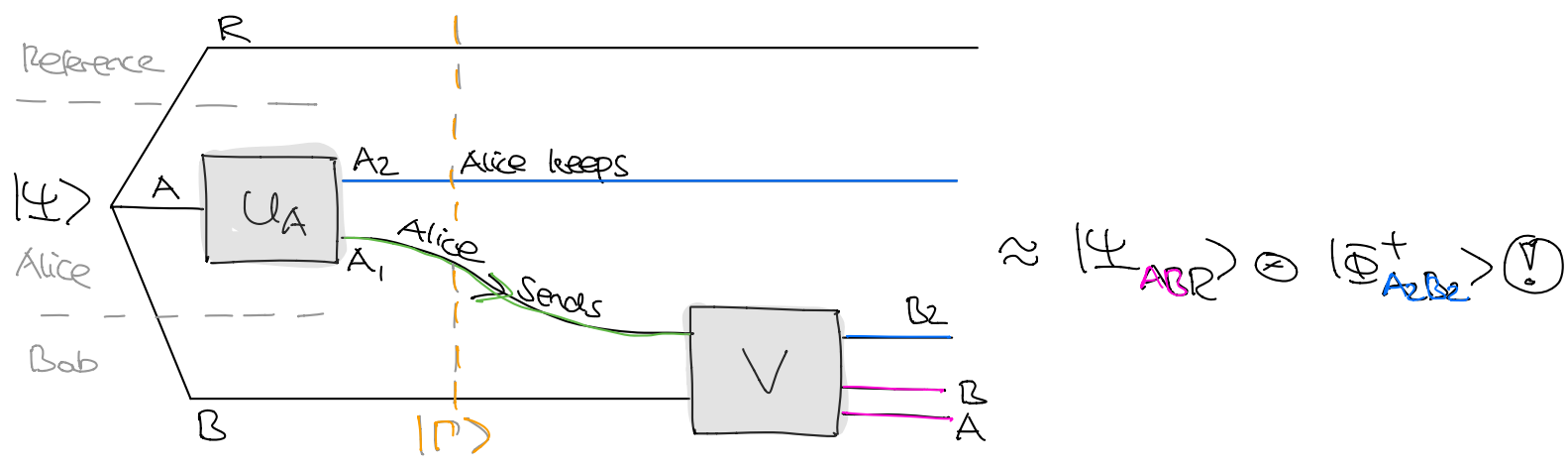
Pf: Send qubits via teleportation. (CI)

Exercises: Noisy teleportation

How to solve it?

Consider first $|\psi_{ABR}\rangle^{\otimes n}$ no $|\psi_{ABR}\rangle$ general pure state.

WANT: Unitary U_A and isometry V s.th.



and A_2 as large as possible and A_1 as small as possible.

Decoupling approach: Enough to consider

$$|\Gamma_{ABR}\rangle := (U_A \otimes I_B \otimes I_R) |\psi_{ABR}\rangle$$

* If V exists: $\Gamma_{A_2 R} \approx \frac{I_{A_2}}{d_{A_2}} \otimes \psi_R =: \tilde{\Gamma}_{A_2 R}$ (*)

↑ "decoupled"

* This is also sufficient, i.e. (*) \Rightarrow (v) for suitable V ! Get existence of V for free!

Idea: $\tilde{\Gamma}_{A_2 R}$ has purification $|\Gamma_{ABR}\rangle$

$$\tilde{\Gamma}_{A_2 R} \sim |\phi_{A_2 B_2}^+\rangle \otimes |\psi_{ABR}\rangle$$

HW

↑ $F \approx 1$ by (*) $\xrightarrow[\text{Heisenberg}]{\text{Uhlmann}}$ \approx up to isometry $A_1, B \xrightarrow{V} B R B_2$ (†)

How to achieve (*)?

Decoupling theorem: If $\psi_{AR} \in D(A \otimes R)$ where $\psi = \psi_A \otimes \psi_R$:

$$\int dU_A \left\| \text{tr}_{A_1} \left[(U_A^\dagger \otimes I_R) \psi_{AR} (U_A \otimes I_R) \right] - \frac{I_{A_2}}{d_{A_2}} \otimes \psi_R \right\|_1^2$$

$$\leq \frac{d_{AR}}{d_{A_1}^2} \text{tr}[\psi_{AR}^2]$$

← The smaller the more we trace out!

Proof? **HW!**

Let's return to $|\psi_{ABR}\rangle^{\otimes n}$ and solve the state merging problem. How to choose $|\psi\rangle$? Final trick: Typical subspaces! Just like for compression...

Recall: Typical subspaces $S_{n,\epsilon} \subseteq \mathcal{D}^{\otimes n}$ for $g^{\otimes n}$ satisfy

- ① $\text{tr}[\Pi_{n,\epsilon} g^{\otimes n}] \rightarrow 1$ as $n \rightarrow \infty$ where $\Pi_{n,\epsilon}$ orthog. projection
- ② eigenvalues of $\Pi_{n,\epsilon} g^{\otimes n} \Pi_{n,\epsilon}$ in $2^{-n(H(g) \pm \epsilon)}$
- ③ $\dim S_{n,\epsilon} \leq 2^{n(H(g) + \epsilon)}$ and $\geq 2^{n(H(g) - \epsilon)} \text{tr}[\Pi_{n,\epsilon} g^{\otimes n}]$
 $\geq 2^{n(H(g) - 2\epsilon)}$ for large n

Define

$$|\psi_{ABR}\rangle = (\Pi_{n,\epsilon}^A \otimes \Pi_{n,\epsilon}^B \otimes \Pi_{n,\epsilon}^R) |\psi_{ABR}\rangle^{\otimes n} \in \tilde{\mathcal{D}} \otimes \tilde{\mathcal{B}} \otimes \tilde{\mathcal{R}} \subseteq (\mathcal{D} \otimes \mathcal{B} \otimes \mathcal{R})^{\otimes n}$$

where $\Pi_{n,\epsilon}^A$ orthogonal projection onto typical subspace $\tilde{\mathcal{D}} := S_{n,\epsilon}(\psi_A) \subseteq \mathcal{D}^{\otimes n}$

$$\Rightarrow d_{\tilde{\mathcal{D}}} d_{\tilde{\mathcal{R}}} \stackrel{\textcircled{3}}{\leq} 2^{n(H(A) + H(R) + 2\epsilon)} \quad \text{tr}[\psi_{\tilde{\mathcal{B}}}^2]$$

$$\text{tr}[\psi_{\tilde{\mathcal{A}\tilde{\mathcal{R}}}}^2] = \text{tr}[\psi_{\tilde{\mathcal{B}}}^2] \leq \text{tr}[(\Pi_{n,\epsilon}^B \psi_B^{\otimes n} \Pi_{n,\epsilon}^B)^2] \stackrel{\textcircled{2}, \textcircled{3}}{\leq} 2^{-n(H(B) - 3\epsilon)}$$

ESET $\text{tr}_x[(P \otimes I) S_{xy} (P \otimes I)] \leq S_F$ rank $\leq 2^{n(H(B) + \epsilon)}$
eigenvalues $\leq 2^{-n(H(B) - \epsilon)}$

$$\hookrightarrow \text{RHS in Decoupling Thm} \leq \frac{2^{n(I(A:R) + 5\epsilon)}}{d_{\tilde{\mathcal{A}}}^2}$$

RESULT: There exists U_A s.t. $(*)$ holds if we choose $d_{\tilde{\mathcal{A}}} = 2^{\frac{n(I(A:R)}{2} + 3\epsilon)}$

i.e. send qubits at rate $\approx \frac{I(A:R)}{2}$ ☺

And in this case also obtain ebits at rate $\approx \frac{I(A:B)}{2}$ ☺ □

Pf: $\frac{1}{n} \log d_{\tilde{\mathcal{A}_2}} = \frac{1}{n} \log d_{\tilde{\mathcal{A}}} - \left(\frac{I(A:R)}{2} + 3\epsilon \right)$

$$\stackrel{\textcircled{3}}{\Rightarrow} H(A) - \frac{I(A:R)}{2} - 5\epsilon = \frac{I(A:B)}{2} - 5\epsilon \quad \square$$

What we did NOT cover:

Noisy q. Communication channels + their Capacities
to send bits, qubits,...

GOOD LUCK FOR THE EXAM 😊