

This course:

# Intro to QIT from the perspective of symmetries

interdisciplinary field

goal is to leverage

QIT to process information

Storage, transmission, processing

of Q. information; design of high-precision measurements;

Q. cryptography, computation

language/toolbox

"Correlation", "qubit",

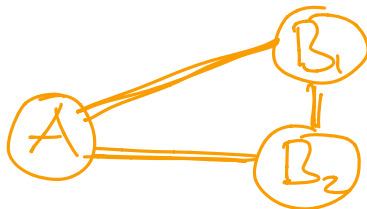
"information", "entropy",

"Complexity", "computation", ...

QIT is "synthetic",  
Why symmetries?

$|4\rangle^{\otimes n}$

many copies  $\rightarrow$  perm. sym.



monogamy of entanglement

$S(\rho) = S(U\rho U^\dagger)$  entropy  $\rightarrow$  unitary sym.

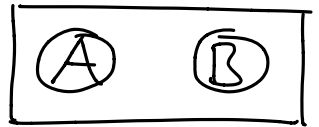
on the technical level, symmetries enter through:

Commut. relations, group actions,  
representation theory

# Introduction to Quantum Mechanics

Axioms are a 1<sup>st</sup> attempt - we will improve them soon!

Axiom (A): To every q. system, we associate a **Hilbert Space**  $\mathcal{H}$ . For a joint system composed of subsystems with Hilbert spaces  $\mathcal{H}_A, \mathcal{H}_B$  the Hilbert space is  $\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B$ .



\* fin-dim HS = vector space & inner product  $\langle - | - \rangle$

\* Throughout this course:  $\dim \mathcal{H} < \infty$

\* Simplest system: **qubit**  $\mathbb{C}^2$

↳ n qubits:  $\mathbb{C}^2 \otimes \dots \otimes \mathbb{C}^2 = (\mathbb{C}^2)^{\otimes n}$   $\dim = 2^n$

Axiom (B): **Unit** vectors  $|\psi\rangle \in \mathcal{H}$  describe the **state** of a system.

\* Dirac notation:  $|\psi\rangle = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \vdots \end{pmatrix}$  "ket"

$\langle\psi| = |\psi\rangle^\dagger = (\bar{\psi}_1 \ \bar{\psi}_2 \ \dots)$  "bra"

$X^\dagger = \bar{X}^T$   
adjoint

$\langle\psi|\phi\rangle = \sum_i \bar{\psi}_i \phi_i$  "bra-ket"

$\langle\psi|\psi\rangle = \|\psi\|^2$

anti-linear in 1<sup>st</sup> argument !!!

\* Standard ("computational") basis:

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad \text{of } \mathbb{C}^2$$

A  
BIT

$$|i_1, i_2, \dots, i_n\rangle = |i_1\rangle \otimes |i_2\rangle \otimes \dots \otimes |i_n\rangle \quad \text{of } (\mathbb{C}^2)^{\otimes n}$$

\* **NOT** every state is a tensor product !!!

e.g.  $\mathbb{C}^2 \otimes \mathbb{C}^2$ :

PSET

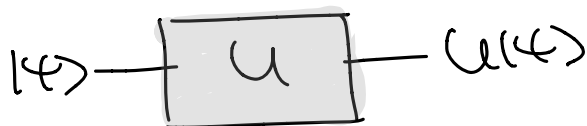
$$|\Phi^+\rangle := \frac{1}{\sqrt{2}} (|0,0\rangle + |1,1\rangle) \neq |\psi\rangle \otimes |\phi\rangle$$

maximally entangled state, ebit, EPR pair

Def:  $|\psi_{AB}\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B$  is called **entangled** if

$$|\psi_{AB}\rangle \neq |\psi_A\rangle \otimes |\phi_B\rangle \quad \text{for all } |\psi_A\rangle \in \mathcal{H}_A, |\phi_B\rangle \in \mathcal{H}_B$$

Axiom (C): Given a **unitary** matrix  $U$  on  $\mathcal{H}$ , the transformation  $|\psi\rangle \mapsto U|\psi\rangle$  is in principle physical.



\* Schrödinger eqn?  
 $-i\partial_t |\psi\rangle = H|\psi\rangle$

\* unitary means  $UU^\dagger = U^\dagger U = I$  (identity matrix)

$$\|U|\psi\rangle\|^2 = \langle \psi | U^\dagger U | \psi \rangle = \langle \psi | \psi \rangle = \|\psi\|^2$$

↳ unit vectors (states) are preserved

Axiom (D): Any Hermitian operator  $X$  on  $\mathcal{H}$  corresponds to an observable quantity: Let  $X = \sum_x x \cdot P_x$  be the **spectral decomposition**. **Born's rule**:

When system is in state  $\psi$ ,

$$Pr(\text{outcome } x) = \langle \psi | P_x | \psi \rangle$$

eigenvalues

projector onto eigenspaces

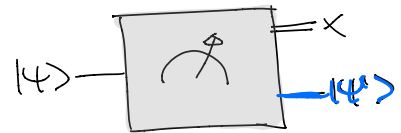
Thus:

$$E[\text{outcome}] = \sum_x x \langle \psi | P_x | \psi \rangle = \langle \psi | X | \psi \rangle$$

$P_x = |e_x\rangle\langle e_x|$   
if non-deg.

Moreover, if the outcome is  $x$  then the **post-meas. state** is

$$|\psi\rangle = \frac{P_x |\psi\rangle}{\|P_x |\psi\rangle\|} = \frac{P_x |\psi\rangle}{\sqrt{\langle \psi | P_x | \psi \rangle}}$$



\* measurement "collapses" the state !!!

↳ **fragility of q. info**, "decoherence"...

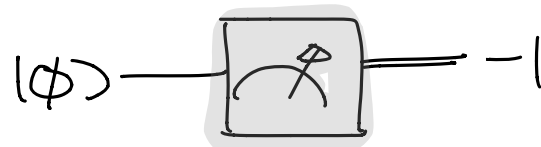
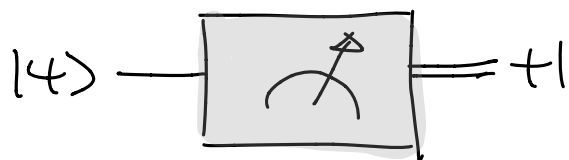
$$P_x |\psi\rangle \perp P_{x'} |\psi\rangle$$

\* **NO collapse**  $\Leftrightarrow |\psi'\rangle = |\psi\rangle \Leftrightarrow P_x |\psi\rangle = \delta_{x,x_0} |\psi\rangle$

$\Leftrightarrow |\psi\rangle$  eigenvector of  $X \Leftrightarrow \langle \psi | P_x | \psi \rangle = \delta_{x,x_0}$

$\Leftrightarrow$  outcome **DETERMINISTIC**

\* When can we distinguish  $|+\rangle, |\phi\rangle$  perfectly?



If orthogonal, i.e.  $\langle +|\phi\rangle = 0$ ! Use observable  $\hookrightarrow$  PSET

$$X = |+\rangle\langle +| - |\phi\rangle\langle \phi| \leftarrow \text{spectral decompo!!!}$$

\*  $|+\rangle$  and  $e^{i\theta}|+\rangle$  are completely indistinguishable

$\hookrightarrow$   $\rho = |+\rangle\langle +|$  density matrix characterizes state up to overall phase

## Measuring a qubit

Infinitely many inequiv. measurements possible. Cf. bit

E.g.

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = |+\rangle\langle +| - |-\rangle\langle -|$$

$$Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = |L\rangle\langle L| - |R\rangle\langle R|$$

$$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = |0\rangle\langle 0| - |1\rangle\langle 1|$$

Spectral decompo  
 $\hookrightarrow$  eigenvalues  $\pm 1$

where  $|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ ,  $|-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$

$|L\rangle = \frac{1}{\sqrt{2}}(|0\rangle + i|1\rangle)$ ,  $|R\rangle = \frac{1}{\sqrt{2}}(|0\rangle - i|1\rangle)$

\*  $\{I, X, Y, Z\}$  is basis of Herm. matrices

\* Pauli matrices do not commute

$$[X, Y] := XY - YX = 2iZ \neq 0 \quad (\text{etc.})$$

↳ Order of measurement matters! = "Joint measurement" does not make sense.

\* Before: Measurement outcome is certain/determined if  $|\psi\rangle$  is eigenvector. BUT: No two Paulis have joint eigenvector.

$$\{X, Z\} := XZ + ZX = 0$$

↳ uncertainty in either  $X$  or  $Z$  measurement (or both)

How to quantify? Consider

$P_X(x)$  = prob of outcome  $x$  when meas.  $X$

$$K_{\psi}(X|\psi) = |P_X(1) - P_X(-1)| = |2P_X(1) - 1| \leq 1$$

\* = 1 if outcome certain,

= 0 if outcome completely random

\* Uncertainty principle: For every state  $|\psi\rangle$ ,

$$|\langle\psi|X|\psi\rangle| + |\langle\psi|Z|\psi\rangle| \leq \sqrt{2} < 2$$

Proof: Let  $s_x, s_z \in \{\pm 1\}$ . Then:

$$\begin{aligned} & s_x \langle\psi|X|\psi\rangle + s_z \langle\psi|Z|\psi\rangle \\ &= \langle\psi| \underbrace{s_x X + s_z Z}_{=: A} |\psi\rangle \stackrel{\text{CS}}{\leq} \|A\| := \sup_{\|\phi\|=1} \|A\phi\| \end{aligned}$$

operator norm

Note:

$$A^\dagger A = A^2 = I + s_x s_z \underbrace{(XZ + ZX)}_{=0} + I = 2 \cdot I$$

$$\Rightarrow \|A\|^2 = \|A^\dagger A\| = 2 \quad \square$$

NB: Works for every state (just like it should)!

↳ Pset for more discussion.

