

Formalism of Q. Information Theory

§2, §3.1 in Watrous

To every (q.) system X , we associate a Hilbert space \mathcal{H} .

* This course: $\dim \mathcal{H} < \infty$

* Dirac notation: $| \psi \rangle \in \mathcal{H}$, $\langle \phi | = (\phi)^* \in \mathcal{H}^*$, $\langle \phi | \psi \rangle$ inner product

$$\mathcal{H} = \mathbb{C}^d: \quad \begin{pmatrix} \psi_1 \\ \vdots \\ \psi_d \end{pmatrix} \quad \begin{matrix} \text{"ket"} \\ \| \\ \langle \phi_1 | \dots | \phi_d \rangle \end{matrix}$$

$$\begin{matrix} \text{adjoint} \\ \downarrow \\ \langle \phi | \psi \rangle \end{matrix} \quad \begin{matrix} \text{"bra"} \\ \| \\ \langle \phi | \end{matrix} \quad \begin{matrix} \text{"bra-ket"} \\ \| \\ \sum_i \phi_i | \psi_i \rangle \end{matrix}$$

anti-linear

EX CLASS

$\| \psi \| = 1: |\psi\rangle \langle \psi| = \text{orthog. projection onto } \mathbb{C}|\psi\rangle$ CLEAR?

$$[(\psi \otimes \psi)^* = \langle \psi^* | \psi \rangle^* = \langle \psi | \psi \rangle; \underbrace{\psi \otimes \psi \otimes \psi \otimes \psi}_{=1} = \psi \otimes \psi, \text{ ran } = \mathbb{C}|\psi\rangle]$$

Register: \sum finite set $\rightsquigarrow \mathcal{H} = \mathbb{C}^\sum$, on B $\{|x\rangle\}_{x \in \sum}$ "standard basis"

e.g. qubit: $\sum = \{0, 1\} \rightsquigarrow \mathcal{H} = \mathbb{C}^2$ w/ standard basis $|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

A (q.) state on X is an element of

$$D(\mathcal{H}) = \{ g \in \text{Pos}(\mathcal{H}), \text{tr}[g] = 1 \}$$

positive semi-definite (PSD) operators

"density operator"

* g pure: $g = |\psi\rangle \langle \psi|$ for some $|\psi\rangle \in \mathcal{H}, \|\psi\| = 1$

* maximally mixed state: $T = \frac{I}{d}$ ($d = \dim \mathcal{H}$)

→ e.s.

* g classical on $\mathcal{H} = \mathbb{C}^\sum$: $g = \sum_x p_x |x\rangle \langle x|$ prob. dist. $g = \begin{pmatrix} p_0 & 0 \\ 0 & p_1 \end{pmatrix}$ qubit

* $D(\mathcal{H})$ is convex: extreme points = pure states. Indeed:

$$\text{SPECTRAL THM: } \forall g \exists \text{ONB } \{|x_i\rangle\}, \text{prob. dist. } \{p_i\}: \underline{g = \sum_i p_i |\psi_i\rangle \langle \psi_i|}$$

What does $D(\mathcal{H})$ look like for a qubit?



Qubits? Pauli matrices $I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$, $X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$, $Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

are basis of 2×2 Herm. matrices (real vector space)

$$\hookrightarrow g = \frac{1}{2} \begin{pmatrix} 1+z & x-iy \\ x+iy & 1-z \end{pmatrix} \text{ Hermitian, } \operatorname{tr} = 1 \quad r = \begin{pmatrix} x \\ y \\ z \end{pmatrix} \text{ Bloch vector}$$

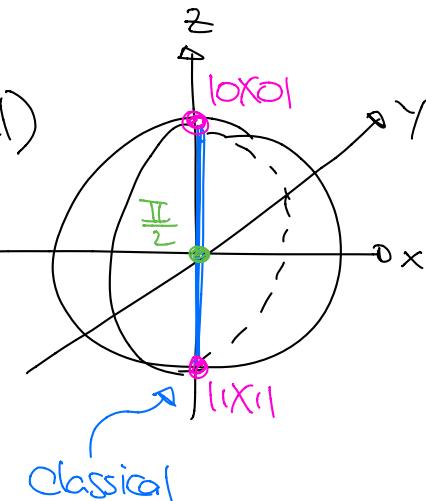
* eigenvalues: $\{p, 1-p\}$ for $p \in \mathbb{R}$

$$\hookrightarrow p(1-p) = \det(g) = \frac{1}{4}(1-x^2-y^2-z^2) = \frac{1}{4}(1-\|r\|^2)$$

* g state: $p(1-p) \geq 0 \iff \|r\| \leq 1$ Bloch ball

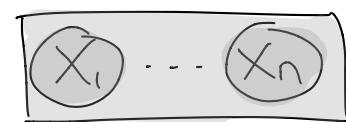
* g pure: $p(1-p) = 0 \iff \|r\| = 1$ Bloch sphere

* g classical: $g = \begin{pmatrix} p & \\ & 1-p \end{pmatrix} \stackrel{\text{Simplex}}{\approx} [0,1] \iff r = \begin{pmatrix} 0 \\ 0 \\ * \end{pmatrix}$



EX. CLASS

For a (g.) system composed of subsystems X_1, \dots, X_n , the Hilbert space is $\mathcal{H} = \mathcal{H}_1 \otimes \dots \otimes \mathcal{H}_n$.



Registers: $\sum = \sum_1 \times \dots \times \sum_n : \mathcal{H} = \mathbb{C}^{\sum} \cong \mathbb{C}^{\sum_1} \otimes \dots \otimes \mathbb{C}^{\sum_n}$
w/ product basis $|x_1, \dots, x_n\rangle = |x_1\rangle \otimes \dots \otimes |x_n\rangle$

e.g. n qubits: $\sum = \{0,1\}^n$, $\mathcal{H} \cong (\mathbb{C}^2)^{\otimes n}$

* product state: $g = g_1 \otimes \dots \otimes g_n$ where $g_i \in D(\mathcal{H}_i)$

* otherwise correlated (e.g. entanglement)
e.g. $\frac{1}{2}(|0,0\rangle\langle 0,0| + |1,1\rangle\langle 1,1|)$

$$= \begin{pmatrix} \frac{1}{2} & 0 & 0 & \frac{1}{2} \\ 0 & \frac{1}{2} & 0 & 0 \\ 0 & 0 & \frac{1}{2} & 0 \\ \frac{1}{2} & 0 & 0 & \frac{1}{2} \end{pmatrix} \in D(\mathbb{C}^2 \otimes \mathbb{C}^2)$$

A measurement (Povm) on X with outcomes in Ω (finite set) is

$$p: \Omega \rightarrow \text{Pos}(\mathcal{H}) \text{ s.t. } \sum_{w \in \Omega} p(w) = I_{\mathcal{H}}$$

"Povm element"

Rule: If measure in state g , obtain outcome

$w \in \Omega$ with probability $p(w) = \operatorname{tr}[p(w)g]$ prob. dist!

$$g \rightarrow \boxed{\text{Povm element}} = w \in \Omega$$

NB: Post-meas. state NOT specified (yet)!

* ρ **projective**: $\rho(\omega)$ projections (\mathcal{H}_ω) \rightarrow your QM class

* **basis measurement**: $\rho(\omega) = |\psi_\omega\rangle\langle\psi_\omega|$ for any $\{|\psi_\omega\rangle\}_{\omega \in \Omega}$
e.g. standard basis of $\mathbb{C}^2 \rightsquigarrow \rho: \Sigma \rightarrow \text{Pos}(\mathcal{H})$

Qubits? Can measure in **standard basis** $\{|0\rangle, |1\rangle\}$, but also in
Hadamard basis $\{|+\rangle, |-\rangle\}$, where $| \pm \rangle = \frac{1}{\sqrt{2}}(|0\rangle \pm |1\rangle)$

e.g. $g = \begin{pmatrix} 2/3 & 1/3 \\ 1/3 & 1/3 \end{pmatrix}$: $P_{\text{std}}(0) = \text{tr}[|0\rangle\langle 0|g] = \langle 0|g|0\rangle = \frac{2}{3}$ } **CANNOT BOTH**
 $P_{\text{had}}(+)=\text{tr}[|+\rangle\langle +|g]=\langle +|g|+\rangle=\frac{1}{2}$ } **BE = 1**

\rightarrow EX CLASS

Discriminating q. states:



WANT: Measurement $\rho: \{0,1\} \rightarrow \text{Pos}(\mathcal{H})$ s.t.

$$P_{\text{success}} = \frac{1}{2} \text{tr}[\rho(0)g_0] + \frac{1}{2} \text{tr}[\rho(1)g_1] \text{ maximal!}$$

e.g. for a qubit:

* $g_0 = |0\rangle\langle 0|$, $g_1 = |1\rangle\langle 1|$: $P_{\text{success}} = 100\%$ via std. basis meas.

* $g_0 = \begin{pmatrix} 2/3 & 0 \\ 0 & 1/3 \end{pmatrix}$, $g_1 = \begin{pmatrix} 1/3 & 0 \\ 0 & 2/3 \end{pmatrix}$: $P_{\text{success}} = \frac{2}{3}$ ————— **optimal?**
in general?

In general: Set $\Omega := \rho(0)$. Then:

$$P_{\text{success}} = \frac{1}{2} \text{tr}[\Omega g_0] + \frac{1}{2} \text{tr}[(I-\Omega)g_1] = \frac{1}{2} + \frac{1}{2} \text{tr}[\Omega(g_0 - g_1)]$$

$$\leq \frac{1}{2} + \frac{1}{4} \|g_0 - g_1\|_1$$

trace distance

Here:

* $\|A\|_1 := \sum_i |\alpha_i|$ trace norm of Hermitian A with eigenvalues $\{\alpha_i\}$.

* for $A = g_0 - g_1 = \sum_i \alpha_i |e_i\rangle\langle e_i|$:

$$\text{tr}[Q_A] = \sum_i a_i \underbrace{\text{tr}[Q|_{\{i\}} \times_{\{i\}}]}_{\in [0,1]} \leq \sum_{i: a_i > 0} a_i \stackrel{\substack{\text{if } a_i > 0 \\ \sum a_i = 0}}{\oplus} \frac{1}{2} \sum |a_i|$$

Equality? $Q = \sum_{i: a_i > 0} |_{\{i\}} \times_{\{i\}}$ (projection!)

Thm (Helstrom): The optimal prob. of success is $\frac{1}{2} + \frac{1}{4} \|g_0 - g_1\|_1$.
It is achieved by a projective measurement.

Q. Channels: $g_{in} \rightarrow \boxed{??}$ $\rightarrow g_{out}$

PICTURES

* $I_X[g] = g$ identity channel

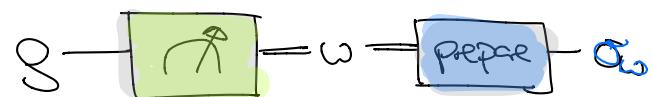
* $\Phi[g] = U g U^\dagger$ for unitary U
"base change"

* $\Phi[g] = \text{tr}[g] \sigma$ for state σ
discards input and prepares state σ

* $\Phi[g] = g \otimes \sigma$ for state σ
prepares additional system in state σ

* "measure and prepare" channel:

home work
:
 $\Phi[g] = \sum_{\omega} \text{tr}[p(\omega) g] \cdot \sigma_{\omega}$



where $p: \Omega \rightarrow \text{Pos}(\omega)$ meas., $\{\omega\}_{\omega \in \Omega}$ states on \mathcal{Y}

* depolarizing channel:

$$\Phi[g] = (1-p) \cdot g + p \cdot \frac{I}{d} \text{tr}[g] \quad \left. \begin{array}{l} \text{replace } g \text{ by } \frac{I}{d} \\ \text{w/ probability } p \end{array} \right\}$$

General math. definition? In a later lecture, we will define the set $C(\mathcal{X}, \mathcal{Y})$

of (q.) channels from system X to system Y .

For today + the homework, only important that:

- ① all examples above are indeed q. channels
- ② every q.channel $\Phi \in C(\mathcal{X}, \mathcal{Y})$ is a linear map $\Phi: L(\mathcal{X}) \rightarrow L(\mathcal{Y})$
linear ops